

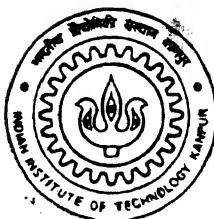
*Editorial*

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# AMMONIA-WATER BINARY-VAPOUR POWER CYCLE - A CHOICE OF FUTURE

by

**LOKESH KUMAR GARG**



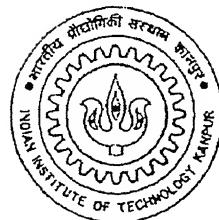
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DEPARTMENT OF MECHANICAL ENGINEERING  
**INDIAN INSTITUTE OF TECHNOLOGY KANPUR**

July, 1999

**AMMONIA-WATER BINARY-VAPOUR  
POWER CYCLE - A CHOICE OF FUTURE**

*A Thesis Submitted*  
in Partial Fulfillment of the Requirements  
for the Degree of  
MASTER OF TECHNOLOGY  
*by*  
LOKESH KUMAR GARG



*to the*  
DEPARTMENT OF MECHANICAL ENGINEERING  
INDIAN INSTITUTE OF TECHNOLOGY KANPUR  
JULY, 1999

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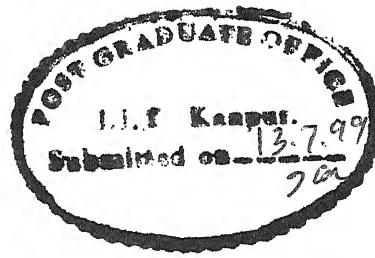
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**Dedicated to**  
My Parents, Sisters and Brothers.

## CERTIFICATE



It is to certify that the work contained in the thesis entitled **Ammonia-Water Binary-Vapour Power Cycle - A Choice of Future**, by **Lokesh Kumar Garg**, has been carried out under my supervision and that this work has not been submitted elsewhere for a degree.

*Chreest* 12/7199

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July, 1999

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# Synopsis

Steam is mostly used in thermal power plants. It has negative condensing pressure below 373 K and high volume flow rate in low-pressure turbines. Binary-vapour cycle is used to overcome these deficiencies.

In the present work analysis of ammonia-water binary-vapour cycle has been carried out by developing thermodynamic functions using least square interpolations and by developing a computer program. Steam and ammonia are used in topping cycle and bottoming cycles respectively, it is shown that the efficiency of cycle depends on pressures of feed water heaters. Pressures of feed water heaters are determined using iterative procedure. Volume flow rate in low-pressure turbine and size of low-pressure turbine are found to be much less in ammonia-water binary-vapour cycle than while using steam only. Mathematical functions for some thermodynamic properties have been developed using least square interpolations. Energy used for evacuation in condenser and evacuation system are trivially eliminated in this cycle. Size of the low-pressure turbines and condenser are reduced significantly compared to existing systems. Diameter of low-pressure turbine in binary-vapour cycle is 20%, 8% at the inlet and exit respectively than while using steam only at 165.4 bar boiler pressure.

# Contents

Dedication	i
Certificate	ii
Acknowledgements	iii
Synopsis	iv
Contents	v
List of Figures	viii
List of Tables	xii
Nomenclature	xiii
<b>1 Introduction</b>	<b>1</b>
1.1 Description . . . . .	1
1.2 Binary-Vapour Cycle . . . . .	3
1.3 Literature Survey . . . . .	5
1.4 Present Work . . . . .	6
<b>2 Binary-Vapour Power System</b>	<b>8</b>
2.1 Binary-vapour System . . . . .	8

2.2	Efficiency of Binary-Vapour Power Cycle . . . . .	10
2.3	Operating Conditions . . . . .	11
<b>3</b>	<b>Mathematical Modelling</b>	<b>13</b>
3.1	Binary-Vapour Power System . . . . .	13
3.2	Steam Power System . . . . .	18
<b>4</b>	<b>Results and Discussions</b>	<b>22</b>
4.1	Optimum Saturation Temperature of Feed Water Heater . . . . .	22
4.2	Efficiency . . . . .	22
4.3	Mass flow Rate in Condenser . . . . .	23
4.4	Mass flow Rate of Ammonia . . . . .	24
4.5	Mass Flow Rate through Boiler . . . . .	24
4.6	Volume Flow Rate of Ammonia . . . . .	25
4.7	Volume Flow Rate in Low-Pressure Stage . . . . .	25
4.8	Low-Pressure Turbine Size . . . . .	26
<b>5</b>	<b>Conclusions and Suggestions</b>	<b>43</b>
5.1	Conclusions . . . . .	43
5.2	Suggestions . . . . .	44
<b>Appendix A</b>		<b>45</b>
<b>A</b>	<b>Chebyshev Polynomials</b>	<b>45</b>
<b>Appendix B</b>		<b>46</b>
<b>B</b>	<b>Thermodynamic Properties of Ammonia</b>	<b>46</b>
B.1	Specific Volume of Saturated Liquid of Ammonia . . . . .	46
B.2	Specific Volume of Saturated Vapour of Ammonia . . . . .	47

B.3	Entropy of Saturated Liquid of Ammonia . . . . .	48
B.4	Entropy of Saturated Vapour of Ammonia . . . . .	48
B.5	Enthalpy of Saturated Liquid of Ammonia . . . . .	49
B.6	Enthalpy of Saturated Vapour of Ammonia . . . . .	50
B.7	Saturation Pressure of Ammonia . . . . .	50
B.8	Enthalpy of Superheated Vapour of Ammonia . . . . .	51
B.9	Entropy of Superheated Vapour of Ammonia . . . . .	52
B.10	Specific Volume of Superheated Vapour of Ammonia . . . . .	55
<b>Appendix C</b>		<b>58</b>
<b>C</b>	<b>Thermodynamic Properties of Steam</b>	<b>58</b>
C.1	Enthalpy of Saturated Liquid . . . . .	58
C.2	Enthalpy of Saturated Vapour . . . . .	59
C.3	Entropy of Saturated Liquid . . . . .	59
C.4	Entropy of Saturated Vapour . . . . .	60
C.5	Specific Volume of Saturated Liquid . . . . .	61
C.6	Specific Volume of Saturated Vapour . . . . .	62
C.7	Saturation Pressure . . . . .	62
C.8	Superheated Properties for the Region $0^{\circ}\text{C} \leq t \leq 350^{\circ}\text{C}$ . . . . .	64
C.9	Superheated Properties for the Region $350^{\circ}\text{C} \leq t \leq 500^{\circ}\text{C}$ . . . . .	65
C.10	Superheated Properties for the Region $500^{\circ}\text{C} \leq t \leq 800^{\circ}\text{C}$ . . . . .	72
<b>Bibliography</b>		<b>73</b>

# List of Figures

1.1	Ammonia-water binary-vapour system	4
2.1	Ammonia-water binary-vapour system	9
3.1	Ammonia-water binary-vapour power cycle	14
3.2	Steam power system	19
3.3	Steam power cycle	20
4.1	Variation in efficiency with saturation temperature of feed water heater of steam cycle with one feed water heater	27
4.2	Variation in efficiency with feed water heater pressure of steam cycle with one feed water heater	28
4.3	Variation in efficiency with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$	28
4.4	Variation in efficiency with boiling pressure of steam at the constant saturation temperature in condenser = $313\text{ K}$	29
4.5	Variation in efficiency of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$	29
4.6	Variation in efficiency of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = $313\text{ K}$	30

4.7 Variation in efficiency of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$ . . . . .	30
4.8 Variation in efficiency of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = $313\text{ K}$ . . . . .	31
4.9 Variation in efficiency of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$ . . . . .	31
4.10 Variation in efficiency of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = $313\text{ K}$ . . . . .	32
4.11 Variation in mass flow rate in condenser of steam cycle with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$ . . . . .	32
4.12 Variation in mass flow rate in condenser with boiling pressure of steam at the constant condenser temperature = $313\text{ K}$ . . . . .	33
4.13 Variation in mass flow rate of ammonia of binary-vapour cycle with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$ . . . . .	33
4.14 Variation in mass flow rate of ammonia of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$ . . . . .	34
4.15 Variation in mass flow rate of ammonia of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = $303\text{ K}$ . . . . .	34

4.16 Variation in mass flow rate of ammonia of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	35
4.17 Variation in mass flow rate in boiler with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	36
4.18 Variation in mass flow rate in boiler with boiling pressure of steam at the constant saturation temperature in condenser = 313 K . . . . .	37
4.19 Variation in mass flow rate in boiler of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	38
4.20 Variation in mass flow rate in boiler of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	38
4.21 Variation in mass flow rate in boiler of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = 303 K . . . . .	39
4.22 Variation in volume flow rate of ammonia with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	39
4.23 Variation in volume flow rate of ammonia of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	40
4.24 Variation in volume flow rate of ammonia of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	40

4.25 Variation in volume flow-rate of ammonia of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser =303 K . . . . .	41
4.26 Variation in volume flow rate at the inlet of low stage turbine with boiling pressure of the steam at the constant saturation temperature in condenser=303 K . . . . . . . . .	41
4.27 Variation in volume flow rate at the inlet of low stage turbine with boiling pressure of the steam at the constant saturation temperature in condenser=303 K . . . . . . . . .	42
4.28 Variation in volume flow rate at the inlet of low stage turbine with boiling pressure of the steam at the constant saturation temperature in condenser=303 K . . . . . . . . .	42

# List of Tables

C.1	Values of $\Delta h$ ( $350^{\circ}C \leq t \leq 500^{\circ}C$ )	69
C.2	Values of $\Delta s$ ( $350^{\circ}C \leq t \leq 500^{\circ}C$ )	70
C.3	Values of $\Delta v$ ( $350^{\circ}C \leq t \leq 500^{\circ}C$ )	71

# Nomenclature

$A$	Area ( $m^2$ )
$c$	Thickness coefficient
$D$	Diameter (m)
$h$	Enthalpy ( $\text{kJ/kg}$ )
$h_f$	Enthalpy of saturated liquid ( $\text{kJ/kg}$ )
$h_g$	Enthalpy of saturated vapour ( $\text{kJ/kg}$ )
$l$	Nozzle height (m)
$\dot{m}$	Mass flow rate ( $\text{kg/sec}$ )
$p$	Pressure (bar)
$p_s$	Saturation pressure (bar)
$p_c$	Critical pressure (bar)
$\dot{Q}$	Volume flow rate ( $\text{m}^3/\text{sec}$ )
$q_c$	Critical liquid density ( $\text{kg/m}^3$ )
$q_l$	Density of saturated liquid ( $\text{kg/m}^3$ )
$R$	Gas constant ( $\text{kJ/kg K}$ )
$s$	Entropy ( $\text{kJ/kg k}$ )
$s_f$	Entropy of saturated liquid ( $\text{kJ/kg k}$ )
$s_g$	Entropy of saturated vapour ( $\text{kJ/kg k}$ )
$T$	Temperature ( K )
$T_c$	Critical temperature ( K )
$t_s$	Saturation temperature ( K )
$v$	Specific volume ( $\text{m}^3/\text{kg}$ )
$v_f$	Specific volume of saturated liquid( $\text{m}^3/\text{kg}$ )
$v_g$	Specific volume of saturated vapour ( $\text{m}^3/\text{kg}$ )

$v_c$	Critical specific volume ( $m^3/kg$ )
$V$	Velocity ( m/sec )
$W_{net}$	Net work ( kW )
$W_p$	Work input to pump ( kW )
$W_t$	Work output of turbine ( kW )
$\alpha$	Nozzle angle
$\eta_c$	Efficiency of the cycle
$\eta_{HE}$	Efficiency of the heat exchanger
$\eta_p$	Pump efficiency
$\eta_t$	Turbine efficiency
$\rho$	Density ( $kg/m^3$ )

# Chapter 1

## Introduction

### 1.1 Description

The ideal power cycle is a Carnot cycle for which the thermal efficiency of heat engine depends on the temperatures  $T_1$  and  $T_2$  of source and sink, respectively.

$$\eta_c = 1 - \frac{T_2}{T_1} \quad (1.1)$$

The processes involved in Carnot cycle are:

- Reversible adiabatic compression
- Reversible heat addition ( isothermal process ) at  $T_1$
- Reversible adiabatic expansion
- Reversible heat rejection ( isothermal process ) at  $T_2$

It is not possible to have actual power cycle working on Carnot principle. The phase change cycle ( Rankine cycle ) approaches the processes of Carnot cycle, i.e., reversible heat addition and reversible heat rejection take place as they occur by phase change processes. Expansion is also reversible adiabatic.

However, for fuel cells the thermal efficiency is not limited by Carnot value. It may have the efficiency of energy conversion far in excess of Carnot efficiency [4].

Only heat addition during sensible heating is not reversible. The working fluid predominately used in phase change cycle is water. It has most of the desirable properties for the use in thermal power cycle as is evident from the following brief description.

1. The *latent heat of vaporization* should be as large as possible and *specific heat* low so that relatively little heat is required to raise the liquid temperature to its boiling point.
2. The *critical temperature* should be well above the metallurgical limit so that reversible heat transfer occurs at constant temperature during phase change of the working fluid at the high temperature of cycle.
3. At the highest operating temperature the pressure should be moderate.
4. The *condenser pressure* should be positive so that there is neither leakage of air into the system nor need for use of the evacuation system.
5. *Saturated vapour line* should lie very close to the path of steam expansion to keep the moisture content at the end of expansion within allowed limit of 10% to 12%.
6. *Specific volume* should be low so that volume flow rate in the turbine is low. Thus, the turbine size becomes small.
7. The *freezing point* should be far below the condensing temperature.
8. It should be *non-flammable, non-poisonous, non toxic*, etc.

9. It should be available in *abundance at low cost*.

Steam satisfies most of the requirements. However, deficiencies with regard to the properties of steam are, negative condensing pressure (0.07384 bar at 313 K), abnormally high boiler pressure (150.0 bar) and tremendously large specific volume ( $19.54 \text{ m}^3/\text{kg}$  at 313 K).

## 1.2 Binary-Vapour Cycle

A typical binary-vapour cycle is shown in figure 1.1. It has high boiling fluid steam in the *topping cycle* and low boiling fluid ammonia in the *bottoming cycle*. The condenser of the topping cycle transfers heat for boiling of working substance of the bottoming cycle to render vapour for the same. The binary-vapour power cycle renders improved thermal efficiency due to use of two fluids. Water-mercury binary-vapour cycle was tried long back [10]. The ammonia-water binary-vapour cycle is a suitable choice to overcome the problems regarding properties mentioned at serial numbers 4, 5, 6 in sec. 1.1.

Despite toxic nature of ammonia, its use has been accepted due to its zero ODP (Ozone depleting potential) and zero GWP (Global warming potential). Ammonia has positive condenser pressure (15.55 bar at 313 K), and very low specific volume ( $0.0833 \text{ m}^3/\text{kg}$  at 313 K). At 313 K ammonia has 1/250 and at 293 K 1/400 specific volume of steam. So steam must be expanded to atmospheric pressure only. Ammonia, as its critical temperature (405.5 K) is low, replaces low-pressure steam turbines for bottoming cycle only. Steam is used in the topping cycle in the ammonia-water binary-vapour cycle. Condenser of topping cycle acts as thermal reservoir for boiling of ammonia of the bottoming cycle. Problem of air leakage into the condenser for ammonia-water binary cycle is eliminated due to positive condenser pressure of ammonia. A

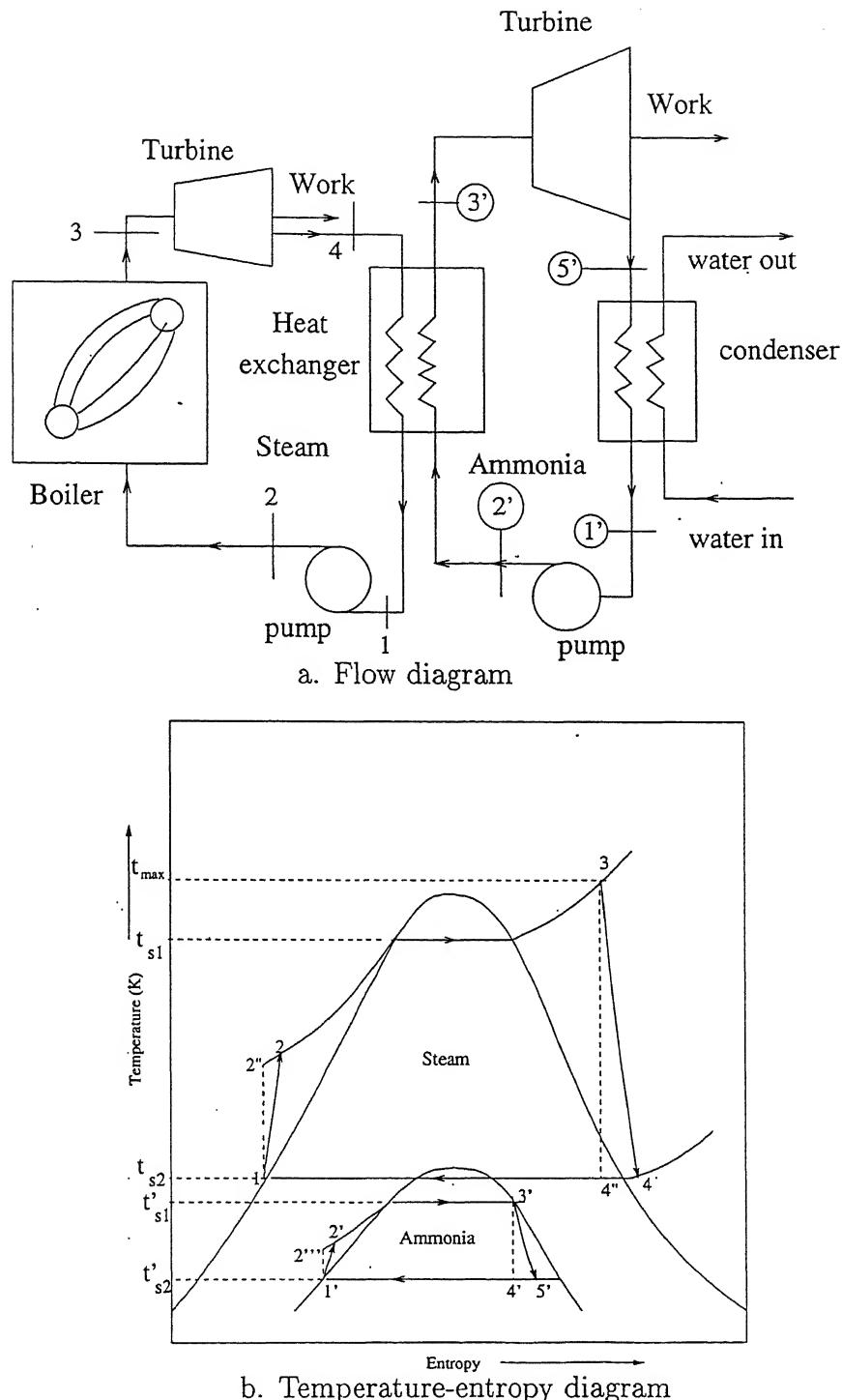


Figure 1.1: Ammonia-water binary-vapour system

continuous evacuation system is trivially eliminated. Further the turbine size of low-pressure stage becomes quite small because of low specific volume of ammonia. A single-flow ammonia turbine can replace several double-flow low-pressure steam turbines [7].

### 1.3 Literature Survey

A high initial temperature is necessary in order to attain a high thermal efficiency for conversion of heat to mechanical energy. Steam suffers from the disadvantages that high pressures and superheating are required in order to obtain high temperature. On the other hand, steam is satisfactory for the utilization of low-pressure, but at low-pressure high specific volume of steam leads to abnormally large size steam turbine. It has therefore been proposed to use top fluids with low-vapour pressure at high temperature, steam for medium and moderately low-pressure and bottom fluids with low specific volume [7].

In binary-vapour cycles, working substances like diphenyl oxide, mercury in topping cycle with steam in bottoming cycle have been used.

Diphenyl oxide has a high freezing point. Thus, danger of freezing in pipeline is serious. Diphenyl oxide has a tendency to decompose above 753 K and has a very high vapour pressure above the 753 K temperature.

Mercury has a critical temperature greater than 1823 K and critical pressure greater than 200 bar. However, if mercury expands in a turbine to 383 K saturation temperature, its corresponding pressure is highly negative (0.000625 bar at 383 K), i.e. high vacuum and a tremendous volume of mercury has to be handled. Mercury is poisonous and highly injurious to health and attacks most of the metals in common use. Each kW of installed capacity requires 1.8 kg of mercury and 14 to 18 kg of mercury vapour has to be condensed to

vaporise 1 kg of water vapours. With mercury being very costly, the cost of plant becomes enormously high. Mercury does not wet boiler and condenser surfaces, so it gives poor heat transfer characteristics.[9]

Fluids like ammonia, methyl chloride and ethyl bromide are the working media in the bottoming cycle with steam being used in the topping cycle. Some CFCs (R-11, R-114, R-113) were also used for the bottoming cycle. CFCs are being phased out due to Ozone hole and GWP problems. For the top fluid used in conjunction with water, a eutectic mixture of diphenyl and diphenyl oxide has been proposed. The mercury / benzene binary-vapour cycle, suggested for possible adoption, would eliminate water as a working fluid. The low specific heat of liquid benzene, about 1.7 kJ/kg K should be viewed in the light of a comparatively low value of the energy of evaporation. Liquid metals, particularly potassium and sodium oxide are used in nuclear electric power plants devised for space applications. In a potassium/steam binary-vapour cycle, potassium turbine would be supplied with potassium vapour at about 2 bar and 1103 K.

[10]

## 1.4 Present Work

A computer programme has been developed to analyse the water-ammonia binary-vapour cycle with different number of feed water heaters. Saturation temperatures in feed water heaters are determined in a way to get maximum efficiency of the cycle. Saturation temperature in condenser is taken as 303 K, 313 K. Range of saturation temperature in boiler is from 473 K to 623 K. Calculations have been carried out for 100 MW power plant. Maximum temperature in the cycle is taken as 823 K. Ammonia boiling temperature range is from 363 K to 391 K. Volume flow rate at the inlet of low-pressure

turbine in binary-vapour cycle is much less than while using steam only.

# Chapter 2

## Binary-Vapour Power System

### 2.1 Binary-vapour System

Figure 2.1 shows a simple binary-vapour system. In this system two working fluids are selected on the basis of their desirable properties, such that the overall efficiency of the system gets improved compared to steam power plant cycle. In topping cycle steam is used with ammonia in bottoming cycle. Heat rejected in the condensation of steam is used for boiling of ammonia. Water from pump 1 at point 6 enters in the boiler. Heat is added to water, to get superheated steam. At point 7, temperature of steam is 823 K. The superheated steam at state 7 enters the turbine, where it expands to the pressure  $p_2$ , expansion process in turbine is not isentropic. After expansion  $m_1$  mass of steam is extracted for feed water heating in feed water heater FH 1. Remaining mass ( $m - m_1$ ) enters the turbine at state point 9, where it expands to pressure  $p_3$ . After expansion  $m_2$  mass of steam is extracted for feed water heating in feed water heater FH 2, remaining mass ( $m - m_1 - m_2$ ) enters the turbine at state point 11, where it expands to a pressure  $p_4$ . At state point 13, steam enters in the heat exchanger of steam and ammonia, where steam is condensed by evaporating ammonia.

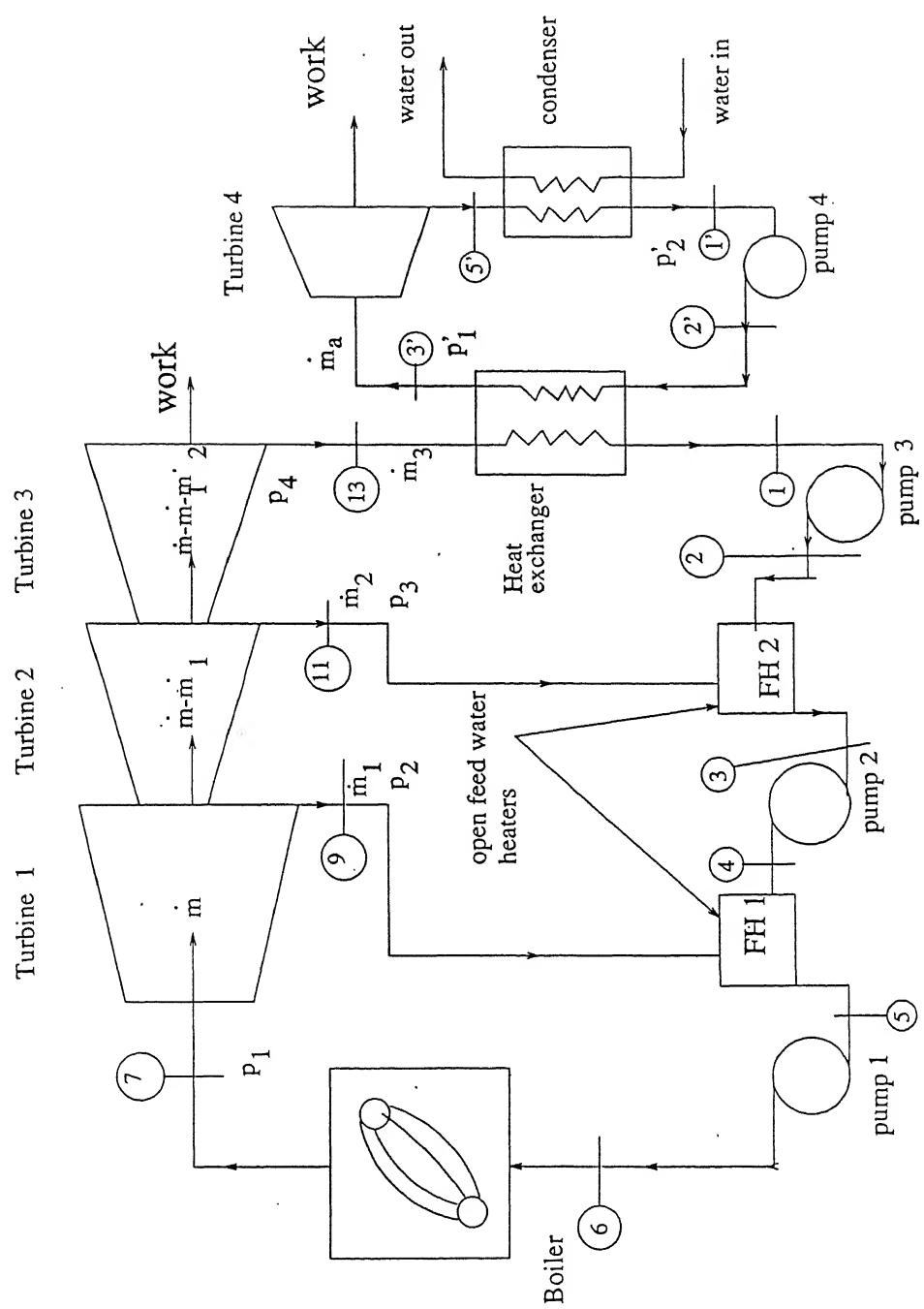


Figure 2.1: Ammonia-water binary-vapour system

Condition of steam at the exit of heat exchanger is saturated liquid. The liquid enters the pump at state point 1 at the condenser pressure  $p_4$ . The liquid is raised upto a pressure  $p_3$ . This mass  $\dot{m}_3$  at pressure  $p_3$  is mixed in feed water heater FH 2 with steam vapour of mass  $\dot{m}_2$  extracted at point 11. Condition of steam at the exit of feed water heater is saturated liquid. At point 3, this mass  $(\dot{m}_2 + \dot{m}_3)$  enters the pump at pressure  $p_3$ . The liquid is raised upto a pressure of  $p_2$ . This mass  $(\dot{m}_2 + \dot{m}_3)$  is mixed in feed water heater FH 1 with steam vapours of mass  $\dot{m}_1$  extracted at point 9. After mixing condition of steam is saturated liquid. At point 5, this mass  $\dot{m}$  enters the pump at pressure  $p_2$ , where it is raised upto a pressure  $p_1$ .

In the bottoming cycle ammonia enters the heat exchanger at pressure  $p_1'$  where it is evaporated by condensing steam. Ammonia is evaporated till it becomes saturated vapour. At state point 3' it enters the turbine, where it expands to pressure  $p_2'$  (condenser pressure). Ammonia at point 5' enters the condenser where its latent heat is removed by circulating water. At point 1' ammonia enters the pump at pressure  $p_2'$  where its pressure is raised upto pressure  $p_1'$ . At point 2' ammonia enters the heat exchanger.

## 2.2 Efficiency of Binary-Vapour Power Cycle

Let  $\eta_1$  and  $\eta_2$  be efficiencies of topping cycle and bottoming cycle respectively, then the binary-vapour cycle efficiency ( $\eta$ ) is given as [8] :

$$\eta = \eta_1 + \eta_2 - \eta_1\eta_2 \quad (2.1)$$

Equation 2.1 can be written as

$$\eta = \eta_1 + \eta_2(1 - \eta_1) \quad (2.2)$$

$$\eta = \eta_2 + \eta_1(1 - \eta_2) \quad (2.3)$$

$\eta_1 < 1, \eta_2 < 1$ , so  $\eta > \eta_1, \eta > \eta_2$

If  $\eta_1 = 0.3, \eta_2 = 0.3$  then  $\eta = 0.51$

So binary-vapour cycle efficiency is higher than cycle using single fluid.

## 2.3 Operating Conditions

The operating conditions for the present system have been decided based on pressures and temperatures of existing power plants. For ammonia side the upper temperature has been decided based on the positive condensing pressure of low-pressure steam turbine. From this consideration condensation temperature of steam is kept 10 K above the ammonia boiling temperature. Ammonia has a critical temperature of 405.5 K. Saturation pressure of steam is 1.01325 bar at 373 K. So we keep minimum condensation temperature of steam as 373 K to avoid vacuum problem, so minimum boiling temperature of ammonia is (373 K-10 K), i.e., 363 K. We also fix the maximum limit of boiling temperature of ammonia as 391 K well below the critical temperature (405.5 K) of ammonia. Critical temperature of steam is 647 K. Maximum saturation temperature of steam in the boiler is taken as 623 K. At 623 K saturation pressure of steam is 165.4 bar. Steam is superheated upto 823 K. Feed water heater pressures  $p_2, p_3$  are determined in a way to get maximum efficiency of cycle. Condensation temperature of ammonia is taken as 303 K, 313 K. At 313 K saturation pressure of ammonia (15.5 bar) is above the atmospheric pressure. So in ammonia-water

binary-vapour cycle there is no problem of vacuum. Heat exchanger efficiency is taken as 90%. Expansion processes in the turbines are not isentropic, isentropic efficiency of turbine is taken as 86%. Pump efficiency is taken as 75%.

# Chapter 3

## Mathematical Modelling

### 3.1 Binary-Vapour Power System

Figure 2.1 shows the schematic diagram of binary-vapour cycle system with two feed water heaters in the topping cycle . This cycle is shown on T-s diagram in figure 3.1 .

$\dot{m}$ =mass flow rate in the boiler (kg/sec)

$\dot{m}_1$ =mass flow rate in the high pressure feed water FH 1 (kg/sec)

$\dot{m}_2$ =mass flow rate of in the low-pressure feed water heater FH 2 (kg/sec)

$\dot{m}_3$ =mass flow rate of steam in the heat exchanger (kg/sec)

$\dot{m}_a$ =mass flow rate of ammonia (kg/sec)

$p_1$ =boiler pressure (bar)

$p_2$ =saturation pressure in high pressure feed water heater FH 1 (bar)

$p_3$ =saturation pressure in low-pressure feed water heater FH 2 (bar)

$p_4$ =steam condensing pressure (bar)

$p'_1$ =ammonia boiling pressure (bar)

$p'_2$ =condenser pressure (bar)

$t_7$  =temperature at the exit of boiler=823 K

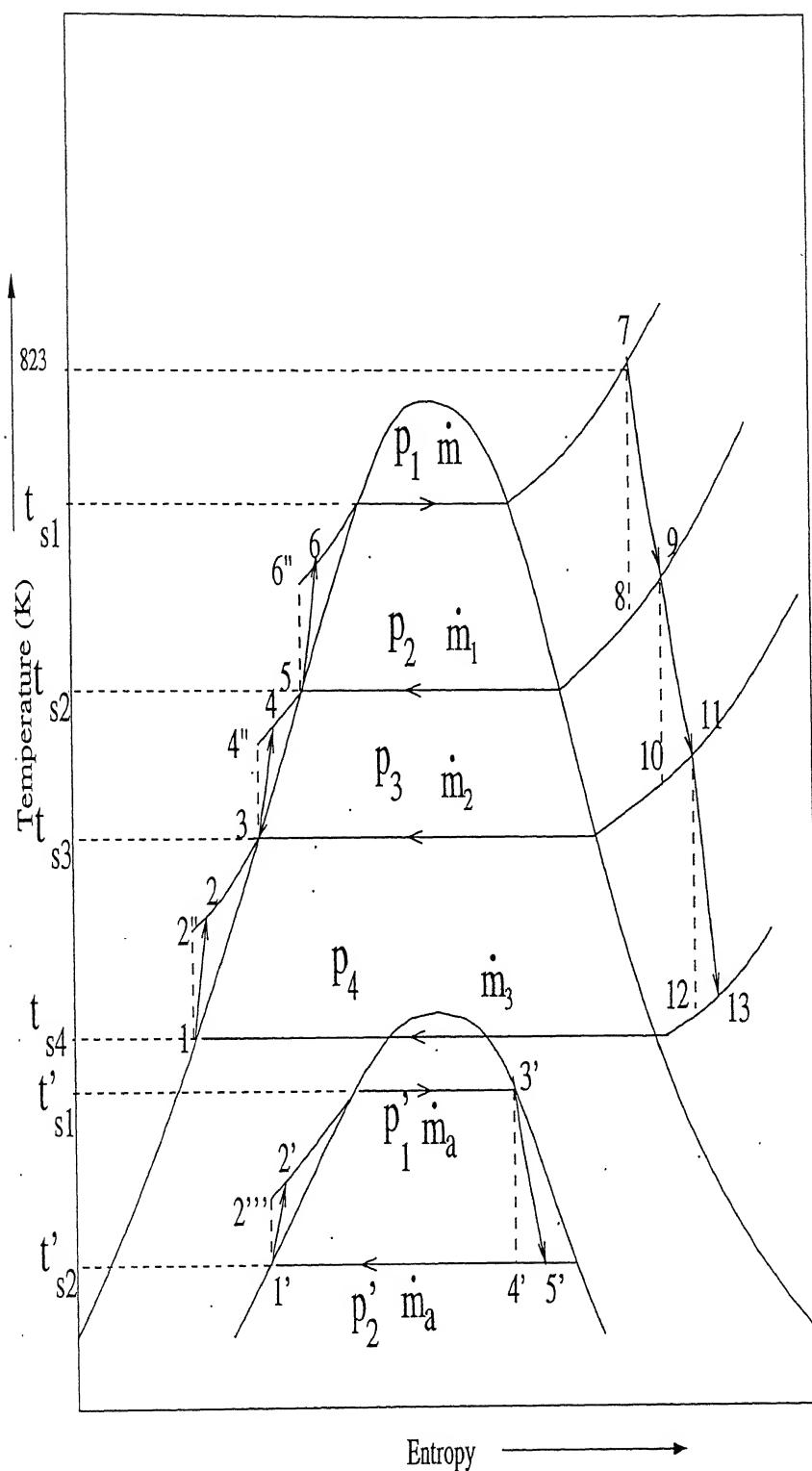


Figure 3.1: Ammonia-water binary-vapour power cycle

$t_{s1}, -t_{s2}, t_{s3}, t_{s4}$  are saturation temperature of the steam at the pressure  $p_1, p_2, p_3, p_4$  respectively.  $t'_{s1}, t'_{s2}$  are saturation temperature of ammonia at the pressures  $p'_1, p'_2$  respectively. Expansion processes in the turbine (7-9, 9-11, 11-13, 3' - 4') are not isentropic. Range of  $p_1$  is from 15 bar to 165.4 bar. A temperature difference of 10 K is maintained between steam condensing temperature and ammonia boiling temperature. Condition of ammonia at point 3' is saturated vapour.

$\eta_{HE}$ =efficiency of the heat exchanger between steam and ammonia=0.90

$\eta_t$ =isentropic efficiency of the turbine =0.86

$\eta_p$ =isentropic efficiency of the pump=0.75

$t'_{s1}=363$  K, 373 K, 383 K, 391 K

$t'_{s2}=303$  K, 313 K

$$t_{s4} = t'_{s1} + 10 \quad (3.1)$$

$$h_{2''} = h_1 + v_1(p_3 - p_4) \quad (3.2)$$

$$h_2 = h_1 + \frac{(h_{2''} - h_1)}{\eta_p} \quad (3.3)$$

$$h_{4''} = h_3 + v_3(p_2 - p_3) \quad (3.4)$$

$$h_4 = h_3 + \frac{(h_{4''} - h_3)}{\eta_p} \quad (3.5)$$

$$h_{6''} = h_5 + v_5(p_1 - p_2) \quad (3.6)$$

$$h_6 = h_5 + \frac{(h_{6''} - h_5)}{\eta_p} \quad (3.7)$$

$$s_8 = s_7 \quad (3.8)$$

$$h_9 = h_7 - \eta_t(h_7 - h_8) \quad (3.9)$$

$$s_{10} = s_9 \quad (3.10)$$

$$h_{11} = h_9 - \eta_t(h_9 - h_{10}) \quad (3.11)$$

$$s_{12} = s_{11} \quad (3.12)$$

$$h_{13} = h_{11} - \eta_t(h_{11} - h_{12}) \quad (3.13)$$

$$h_{2'''} = h_{1'} + v_{1'}(p_1' - p_2') \quad (3.14)$$

$$h_{2'} = h_{1'} + \frac{(h_{2'''} - h_{1'})}{\eta_p} \quad (3.15)$$

$$s_{4'} = s_{3'} \quad (3.16)$$

$$h_{5'} = h_{3'} - \eta_t(h_{3'} - h_{4'}) \quad (3.17)$$

Heat rejected in the condensation of the steam is used for boiling of ammonia.

Assume  $\dot{m}_3 = 1 \text{ kg/sec}$

$$\text{Heat taken by ammonia} = \text{Heat rejected by steam} \quad (3.18)$$

$$\dot{m}_a(h_{3'} - h_{2'}) = \eta_{HE}\dot{m}_3(h_{13} - h_1) \quad (3.19)$$

$$\dot{m}_a = \frac{\eta_{HE}\dot{m}_3(h_{13} - h_1)}{(h_{3'} - h_{2'})} \quad (3.20)$$

$$\dot{m}_a = \frac{\eta_{HE}(h_{13} - h_1)}{(h_{3'} - h_{2'})} \quad (3.21)$$

$$\dot{m}_2(h_{11} - h_3) = \dot{m}_3(h_3 - h_2) \quad (3.22)$$

$$\dot{m}_2 = \frac{\dot{m}_3(h_3 - h_2)}{(h_{11} - h_3)} \quad (3.23)$$

$$\dot{m}_2 = \frac{(h_3 - h_2)}{(h_{11} - h_3)} \quad (3.24)$$

$$\dot{m}_1(h_9 - h_5) = (\dot{m}_2 + \dot{m}_3)(h_5 - h_4) \quad (3.25)$$

$$\dot{m}_1 = \frac{(\dot{m}_2 + \dot{m}_3)(h_5 - h_4)}{(h_9 - h_5)} \quad (3.26)$$

$$\dot{m}_1 = \frac{(\dot{m}_2 + 1)(h_5 - h_4)}{(h_9 - h_5)} \quad (3.27)$$

From equation 3.24,  $\dot{m}_2$  is known so  $\dot{m}_1$  is known.

$$\dot{m} = \dot{m}_1 + \dot{m}_2 + \dot{m}_3 \quad (3.28)$$

$$\dot{m} = \dot{m}_1 + \dot{m}_2 + 1 \quad (3.29)$$

From equation 3.24, 3.27  $\dot{m}_2$ ,  $\dot{m}_1$  are known so  $\dot{m}$  is known. We have calculated mass flow rate at each point of the system.

$$W_t = \dot{m}(h_7 - h_9) + (\dot{m}_2 + \dot{m}_3)(h_9 - h_{11}) + \dot{m}_3(h_{11} - h_{13}) + \dot{m}_a(h_{3'} - h_{5'}) \quad (3.30)$$

$$W_t = \dot{m}(h_7 - h_9) + (\dot{m}_2 + 1)(h_9 - h_{11}) + (h_{11} - h_{13}) + \dot{m}_a(h_{3'} - h_{5'}) \quad (3.31)$$

$$W_p = \dot{m}_3 v_1(p_3 - p_4) + (\dot{m}_2 + \dot{m}_3)v_3(p_2 - p_3) + \dot{m}v_5(p_1 - p_2) + \dot{m}_a v_{1'}(p_1' - p_2') \quad (3.32)$$

$$W_p = v_1(p_3 - p_4) + (\dot{m}_2 + 1)v_3(p_2 - p_3) + \dot{m}v_5(p_1 - p_2) + \dot{m}_a v_{1'}(p_1' - p_2') \quad (3.33)$$

$$W_{net} = W_t - W_p \quad (3.34)$$

$$heat\ added = \dot{m}(h_7 - h_6) \quad (3.35)$$

$$efficiency\ of\ cycle = \frac{net\ work}{heat\ added} \quad (3.36)$$

$$\eta_c = \frac{W_{net}}{\dot{m}(h_7 - h_6)} \quad (3.37)$$

$W_{net}$  is power output for  $\dot{m}_3 = 1$  kg/sec. We assume the capacity of power plant as 100 MW. For 100 MW power output we have to calculate the mass flow rate at each point of the system.

$$\dot{m}_3 = \frac{100\ MW}{W_{net}} \quad (3.38)$$

From equation 3.20, 3.23, 3.26, 3.28  $\dot{m}_a$ ,  $\dot{m}_2$ ,  $\dot{m}_1$ ,  $\dot{m}$  is known for 100 MW power output.

$$\text{volume flow rate} = \text{specific volume} * \text{mass flow rate} \quad (3.39)$$

From equation 3.39 we can calculate the volume flow rate at the inlet and exit of all the turbines.

Saturation temperature in boiler is varied in the range of 473 K to 623 K at the interval of 10 K. Condenser temperature  $t'_{s2}$  is kept at 303 K, 313 K. Boiling temperature of ammonia  $t'_{s1}$  is kept at 363 K, 373 K, 383 K, 391 K. Range of saturation temperature in the high pressure feed water heater ( $t_{s2}$ ) and low-pressure feed water heater ( $t_{s3}$ ) is from  $t_{s4}$  to  $t_{s1}$ . First we keep  $t_{s3}$  5 K above the  $t_{s4}$  and vary  $t_{s2}$  in the range of  $(t_{s3} + 5)$  to  $(t_{s1} - 5)$  till the efficiency becomes the maximum. Then again increase  $t_{s3}$  by 5 K and vary  $t_{s2}$  in the range of  $(t_{s3} + 5)$  to  $(t_{s1} - 5)$  till the efficiency becomes the maximum. We increase  $t_{s3}$  till  $(t_{s1} - t_{s3})$  is greater than 10 K and vary  $t_{s2}$  in the range of  $(t_{s3} + 5)$  to  $(t_{s1} - 5)$  at an interval of 10 K. In this way we can find saturation temperatures in the feed water heaters for maximum efficiency. Thermodynamic properties of ammonia and steam are given in the appendix. A computer program has been developed to analyse the cycle and calculate the maximum efficiency, mass flow rate and volume flow rate at the different points.

## 3.2 Steam Power System

Figure 3.2 shows the schematic diagram of steam power cycle system with three feed water heaters. This cycle is shown on T-s diagram in figure 3.3.

$\dot{m}$ =mass flow rate in the boiler (kg/sec)

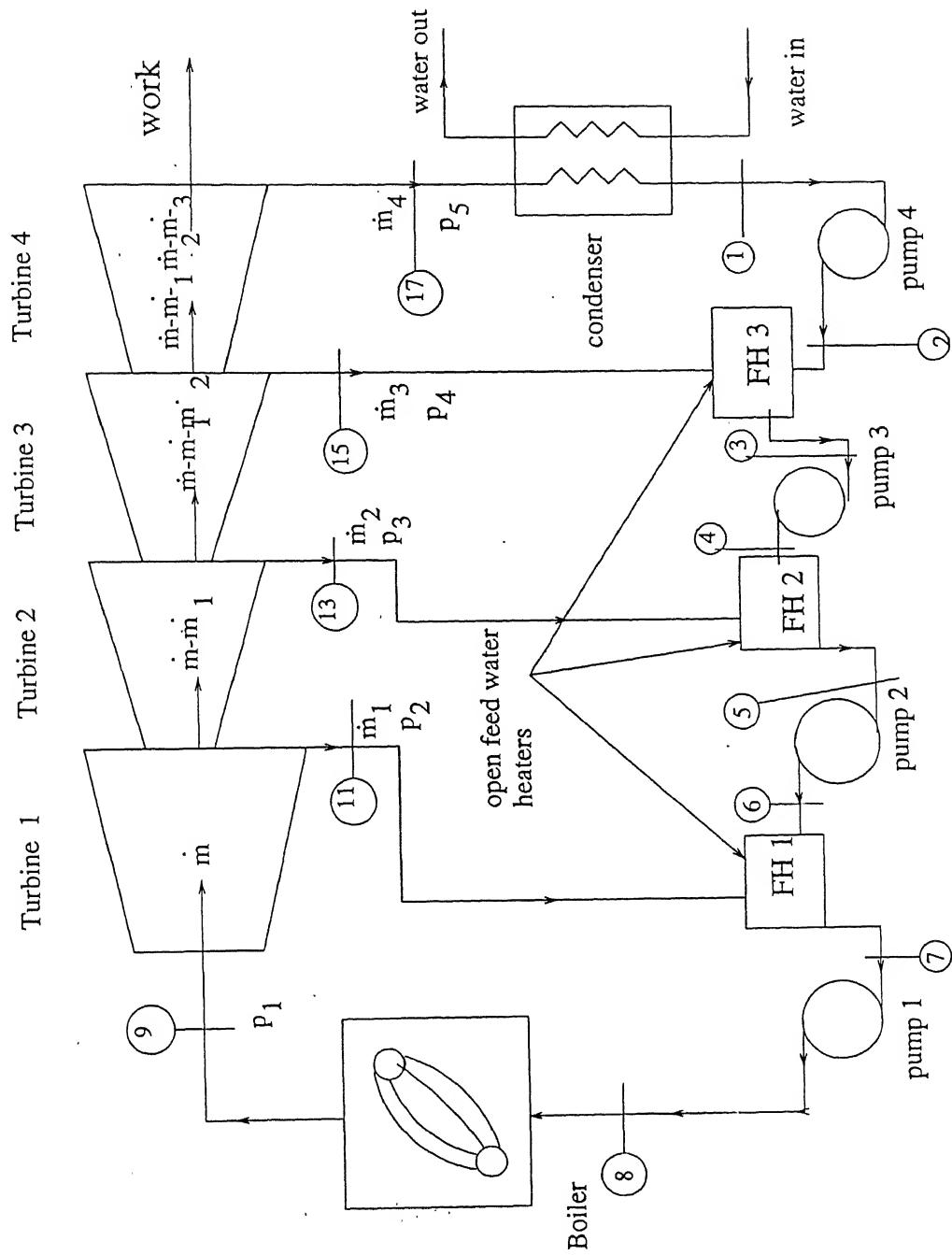


Figure 3.2: Steam power system

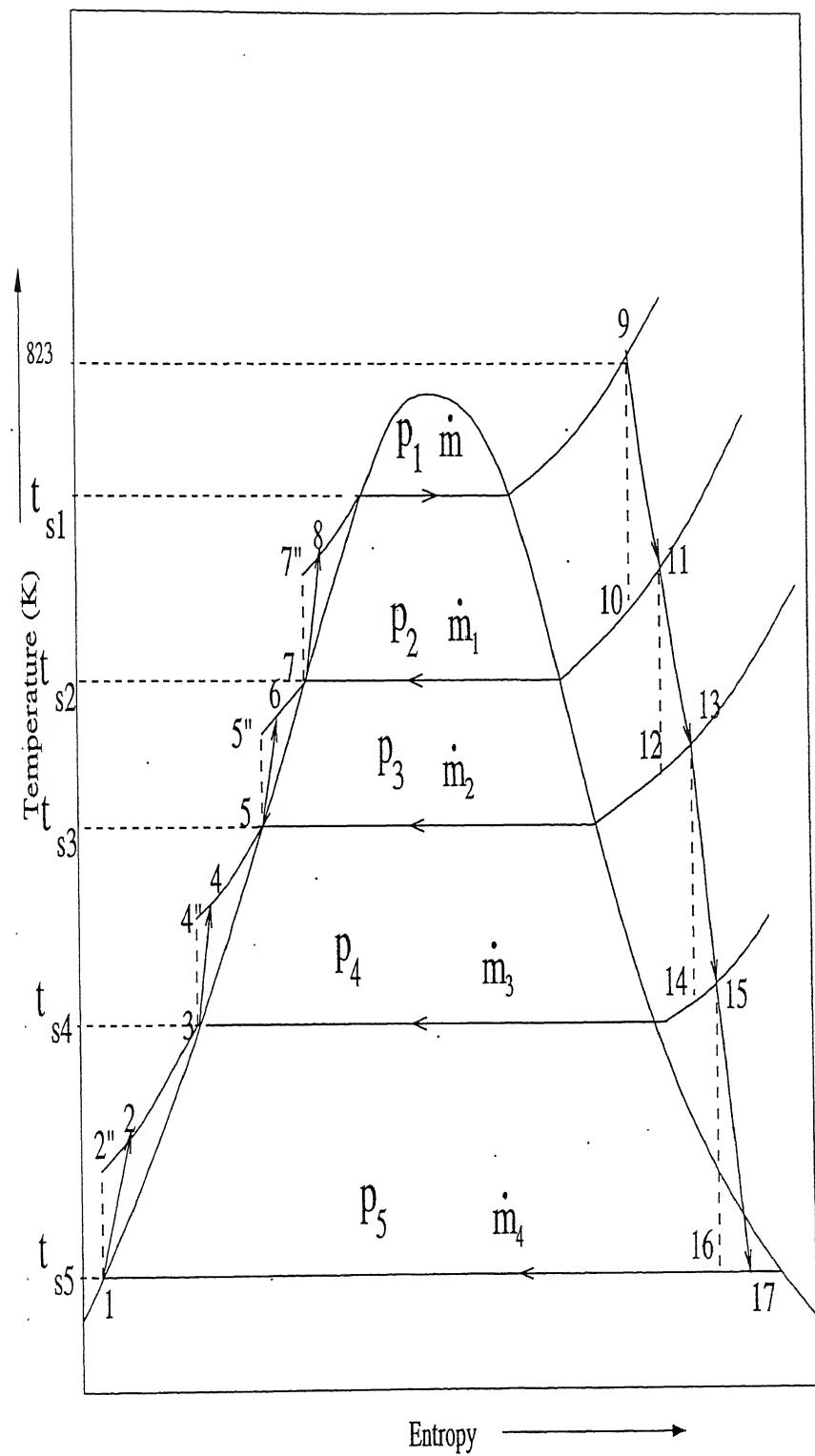


Figure 3.3: Steam power cycle

$\dot{m}_1$ =mass flow rate in feed water heater FH 1 (kg/sec)

$\dot{m}_2$ =mass flow rate of in feed water heater FH 2 (kg/sec)

$\dot{m}_3$ =mass flow rate of in feed water heater FH 3 (kg/sec)

$\dot{m}_4$ =mass flow rate of in the condenser (kg/sec)

$p_1$ =boiler pressure (bar)

$p_2$ =saturation pressure in feed water heater FH 1 (bar)

$p_3$ =saturation pressure in feed water heater FH 2 (bar)

$p_4$ =saturation pressure in feed water heater FH 3 (bar)

$p_5$ =condenser pressure (bar)

$t_9$  =temperature at the exit of boiler=823 K

$t_{s1}, t_{s2}, t_{s3}, t_{s4}, t_{s5}$  are saturation temperature of the steam at the pressure  $p_1, p_2, p_3, p_4, p_5$  respectively. Expansion process in the turbine ( 9-11, 11-13, 13-15, 15-17 ) are not isentropic. Range of  $p_1$  is from 15 bar to 165.4 bar.

$\eta_t$ =isentropic efficiency of the turbine =0.86

$\eta_p$ =isentropic efficiency of the pump=0.75

$t'_{s2}$ =303 K, 313 K

Analysis of this cycle can be done similarly as given in sec. 3.1.

# Chapter 4

## Results and Discussions

### 4.1 Optimum Saturation Temperature of Feed Water Heater

Figure 4.1 shows the variation in efficiency with saturation temperature of feed water heater keeping condenser and boiler pressure constant for steam cycle with one feed water heater. First efficiency increases and becomes maximum before it decreases again. The maximum efficiency is found to occur at the arithmetic mean of saturation temperatures in boiler and condenser [13].

Variation in efficiency with pressure in feed water heater is shown in figure 4.2 keeping condenser and boiler pressure as constant for a steam cycle with one feed water heater.

### 4.2 Efficiency

Figure 4.3, 4.4 shows the variation in maximum efficiency with boiler pressure while saturation temperature of condenser is kept constant. Maximum efficiencies for steam cycle with one, two, three feed water heaters and binary cycle

with one, two feed water heaters has been obtained by using iterative procedure as explained in section no. 3.1. Efficiency at 165.4 bar for steam cycle without feed water heater is 0.38, while for steam cycle with three feed water heaters is 0.45. For the binary-vapour cycle efficiency is slightly less compared to steam cycle because there is a temperature difference of 10 K in heat exchanger and heat exchanger efficiency is 0.9. As number of feed water heaters are increased, efficiency increases.

Figure 4.5, 4.6 shows the variation in maximum efficiency of a binary-vapour cycle without feed water heater for a fixed boiling pressure of steam and different boiling temperatures of ammonia. Efficiency reduces as ammonia boiling temperature is increased.

Figure 4.7, 4.8 shows the variation in maximum efficiency of a binary-vapour cycle with one feed water heater for a fixed boiling pressure of steam and different boiling temperatures of ammonia. Efficiency reduces as ammonia boiling temperature is increased.

Figure 4.9, 4.10 shows the variation in maximum efficiency of a binary-vapour cycle with two feed water heaters for a fixed boiling pressure of steam and different boiling temperatures of ammonia. Efficiency reduces as ammonia boiling temperature is increased.

### 4.3 Mass flow Rate in Condenser

Figure 4.11, 4.12 shows the variation in mass flow rate in condenser with boiling pressure of steam. As number of feed water heaters are increased, condenser mass flow rate reduces.

#### 4.4 Mass flow Rate of Ammonia

Figure 4.13 shows the variation in mass flow rate of ammonia with boiling pressure of steam. As number of feed water heaters are increased, mass flow rate of ammonia reduces. At 165.4 bar boiler pressure and 303 K saturation temperature in condenser, mass flow rate of ammonia is 147 kg/sec for binary-vapour cycle without feed water heater, 123 kg/sec for binary-vapour cycle with two feed water heaters.

Figure 4.14, 4.15 and 4.16 shows the variation in mass flow rate of ammonia with boiling pressure of steam without feed water heater, with one and two feed water heaters respectively at different boiling temperatures of ammonia. Mass flow rate of ammonia decreases as boiling pressure of steam increases. Mass flow rate of ammonia is less for low boiling temperatures of ammonia.

#### 4.5 Mass Flow Rate through Boiler

Figure 4.17, 4.18 shows the variation in mass flow rate in boiler with boiling pressure of steam. For steam cycle without feed water heater and binary-vapour cycle without feed water heater mass flow rate in boiler decreases as boiler pressure increases. Mass flow rate in boiler first decreases, becomes minimum before it increases again for steam cycle with one, two, three feed water heaters and binary-vapour cycle with one, two feed water heaters.

Figure 4.19 shows the variation in mass flow rate in boiler of binary-vapour cycle without feed water heater with boiling pressure of steam. Mass flow rate in boiler increases as ammonia boiling temperature increases.

Figure 4.20, 4.21 shows the variation in mass flow rate in boiler of binary-vapour cycle with boiling pressure of steam with one, two feed water heaters

respectively. Mass flow rate in boiler first decreases, becomes minimum before it increases again.

## 4.6 Volume Flow Rate of Ammonia

Figure 4.22 shows the variation in volume flow rate of ammonia at the inlet of turbine with boiling pressure of steam. As number of feed water heaters are increased, volume flow rate of ammonia decreases. Figure 4.23, 4.24, 4.25 shows the variation in volume flow rate of ammonia at the inlet of turbine with boiling pressure of steam for binary-vapour cycle without feed water heater, with one and two feed water heaters respectively. As ammonia boiling temperature increases, volume flow rate of ammonia decreases.

## 4.7 Volume Flow Rate in Low-Pressure Stage

Figure 4.26 shows the variation in volume flow rate at the inlet of low-pressure stage turbine with boiling pressure of steam for steam cycle with one feed water heater and binary-vapour cycle without feed water heater. Volume flow rate decreases as boiling pressure increases. Volume flow rate for steam cycle with one feed water heater is more than binary-vapour cycle without feed water heater.

Figure 4.27 shows the variation in volume flow rate at the inlet of low-pressure stage turbine with boiling pressure of steam for steam cycle with two feed water heaters and binary-vapour cycle with one feed water heater. Volume flow rate decreases as boiling pressure increases. At 165.5 bar boiler pressure volume flow rate for steam cycle with two feed water heaters is  $34.17m^3/sec$  and  $2.49m^3/sec$  for binary-vapour cycle with one feed water heater.

Figure 4.28 shows the variation in volume flow rate at the inlet of low-pressure stage turbine with boiling pressure of steam for steam cycle with three feed water heaters and binary-vapour cycle with two feed water heaters. Volume flow rate decreases as boiling pressure increases. At 165.4 bar boiler pressure volume flow rate for steam cycle with three feed water heaters is  $68.96\text{m}^3/\text{sec}$  and  $2.92\text{m}^3/\text{sec}$  for binary-vapour cycle with two feed water heaters.

## 4.8 Low-Pressure Turbine Size

At 165.4 bar boiler pressure volume flow rate at the inlet, exit of the low-pressure turbine is  $68.96\text{m}^3/\text{sec}$ ,  $1672\text{m}^3/\text{sec}$  respectively for steam cycle with three feed water heaters and  $2.92\text{m}^3/\text{sec}$ ,  $11.56\text{m}^3/\text{sec}$  respectively for binary-vapour cycle with two feed water heaters. For calculating the size of turbine we assume turbine is of impulse type. [6]

$$A = \pi D l c \sin \alpha \quad (4.1)$$

Where, A=flow area

$\alpha$ =nozzle angle

$l$ =nozzle height

$c$ =thickness coefficient

$D$ =mean diameter of the turbine blade

$$\dot{Q} = A V \quad (4.2)$$

Where,  $\dot{Q}$ =volume flow rate

$V$ =flow velocity

We take  $V=120 \text{ m/sec}$ ,  $150 \text{ m/sec}$  at the inlet, exit of the turbine respectively.

tively.

$$m = \frac{l}{D} \quad (4.3)$$

we take  $m=0.18$ ,  $\alpha = 20^\circ$ ,  $c=0.85$

From equation 4.1, 4.2, 4.3 we get  $D=1.87$  m, 8.23 m at the inlet, exit of the low stage turbine for steam cycle with 3 feed water heaters at 165.4 bar boiler pressure and  $D=0.385$  m, 0.685 m at the inlet, exit of the low-pressure turbine for binary-vapour cycle with 2 feed water heaters at 165.4 boiler pressure. Low-pressure turbine size for steam cycle with 3 feed water heaters is 5, 12 times of binary-vapour cycle with 2 feed water heaters.

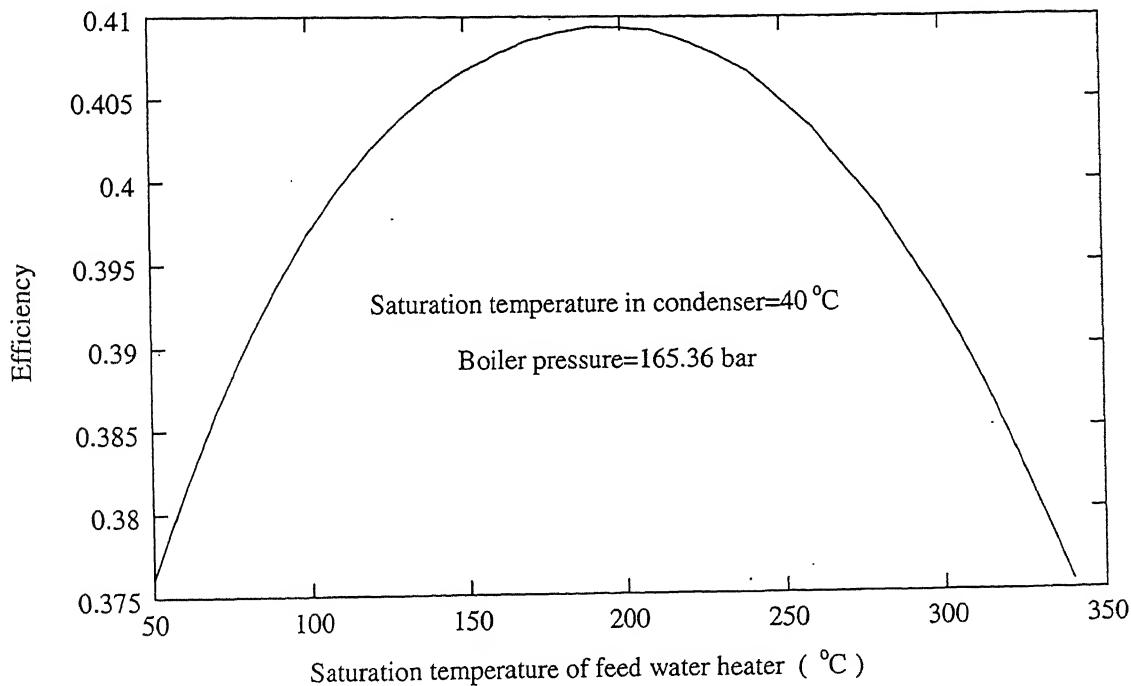


Figure 4.1: Variation in efficiency with saturation temperature of feed water heater of steam cycle with one feed water heater

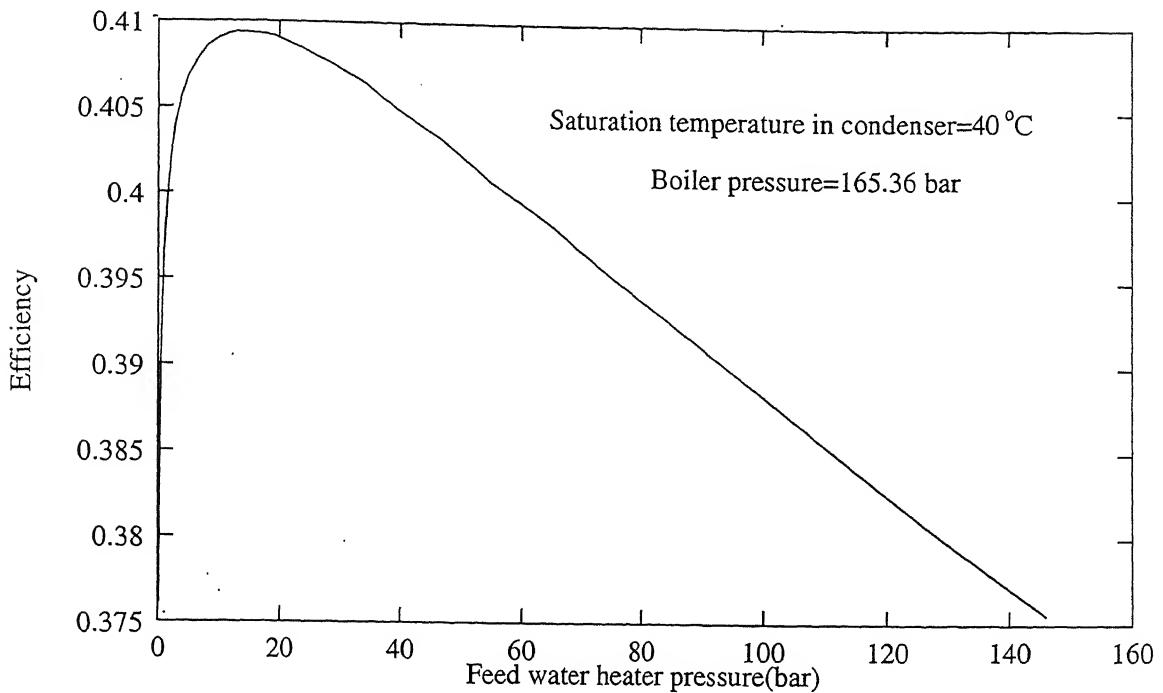


Figure 4.2: Variation in efficiency with feed water heater pressure of steam cycle with one feed water heater

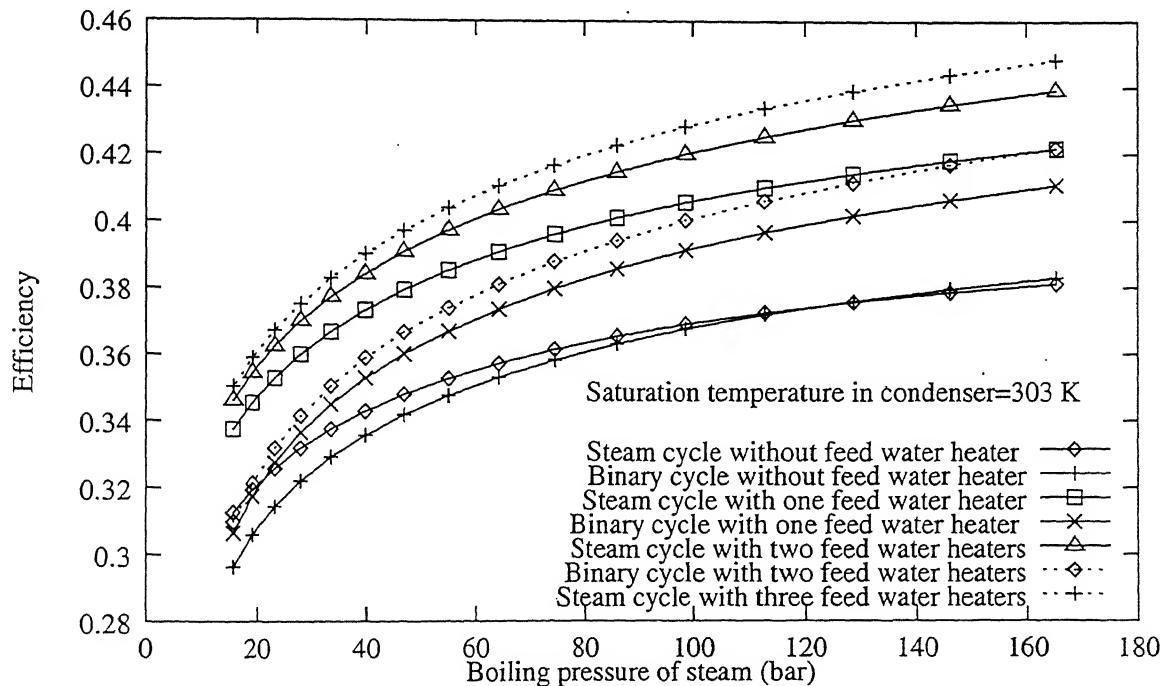


Figure 4.3: Variation in efficiency with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

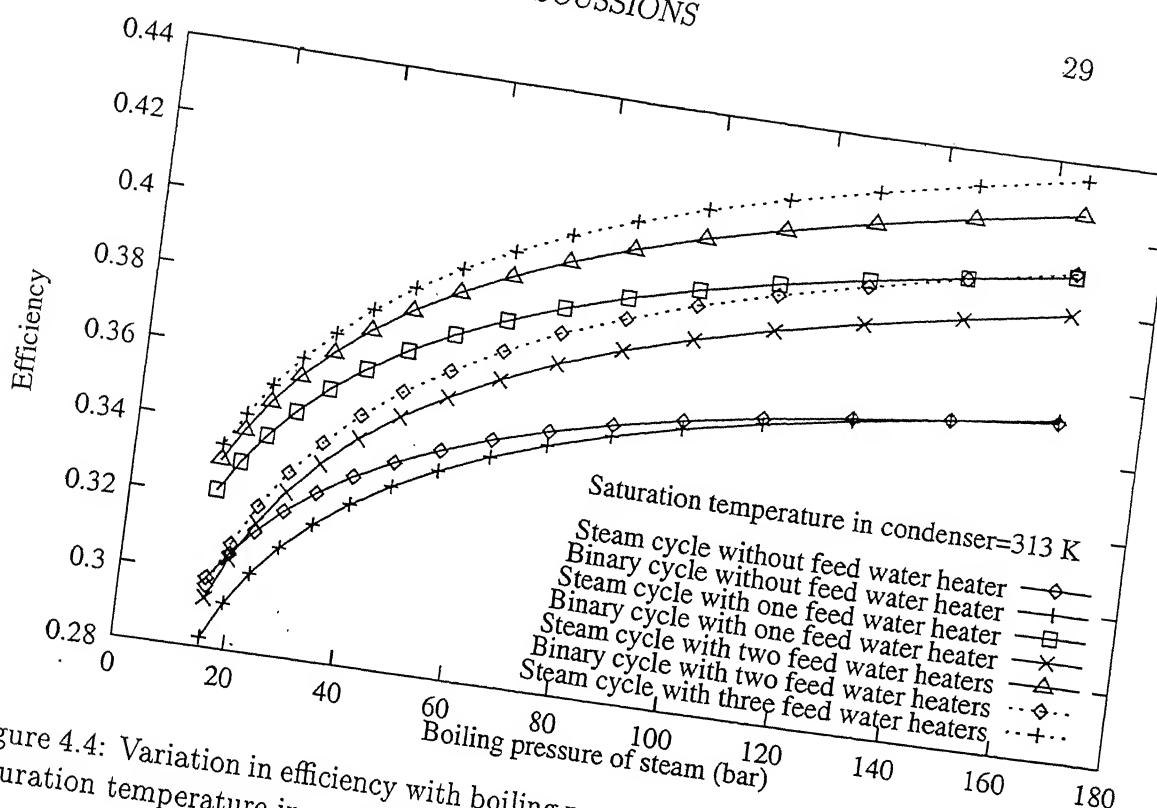


Figure 4.4: Variation in efficiency with boiling pressure of steam at the constant saturation temperature in condenser = 313 K

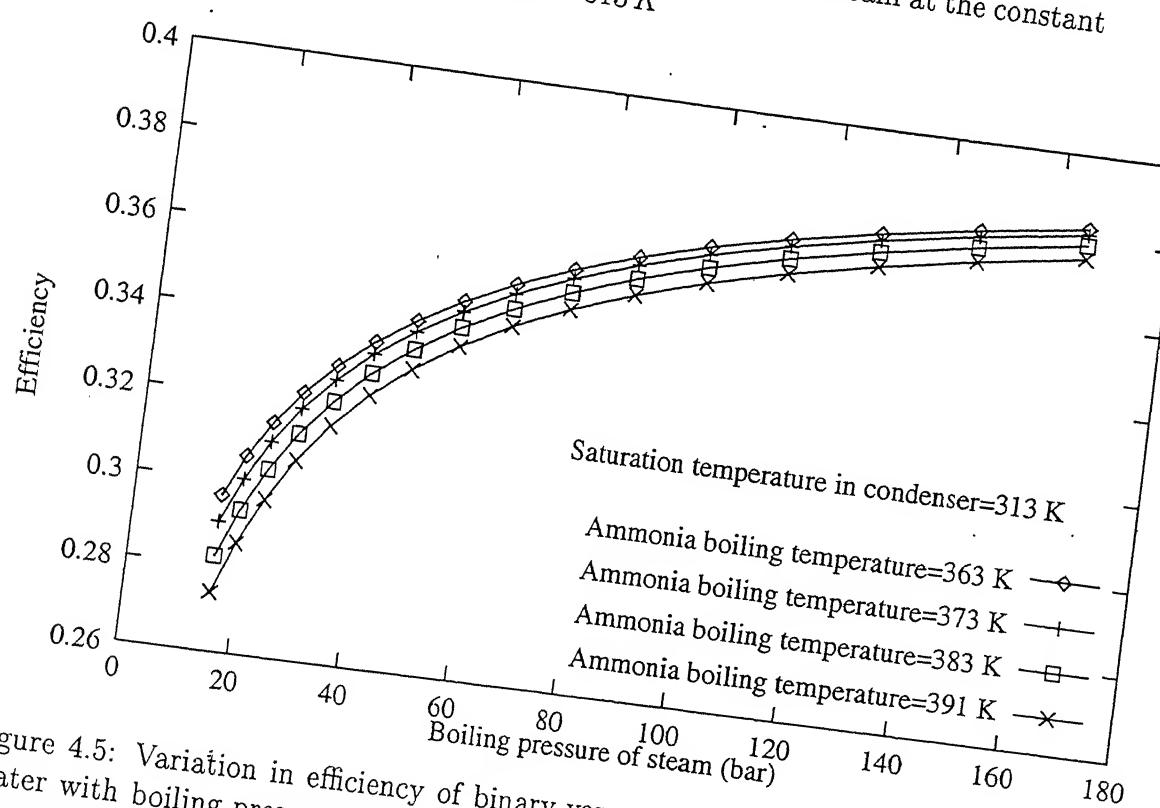


Figure 4.5: Variation in efficiency of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

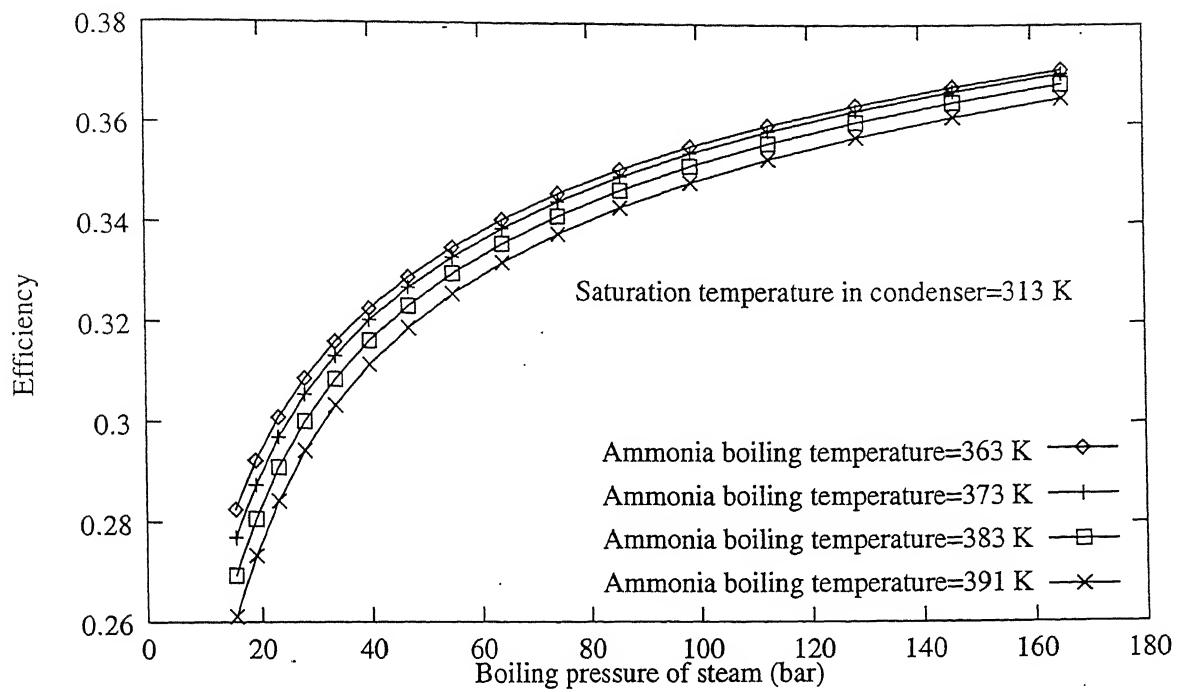


Figure 4.6: Variation in efficiency of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser =313 K

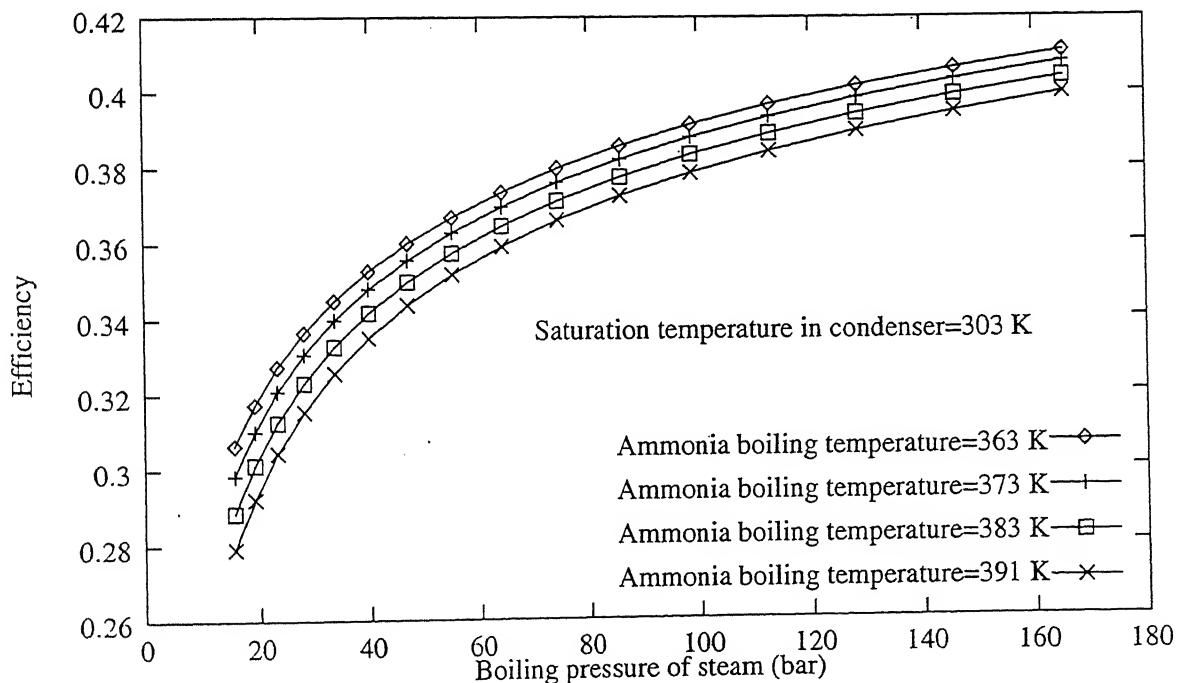


Figure 4.7: Variation in efficiency of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser =303 K

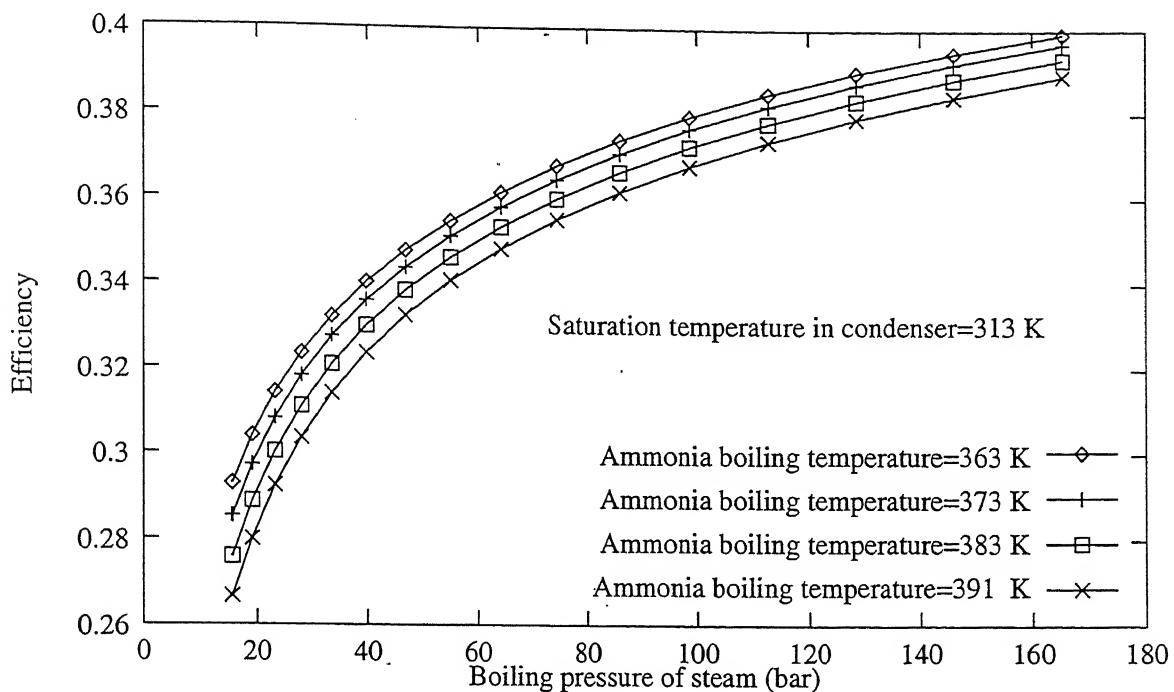


Figure 4.8: Variation in efficiency of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 313 K

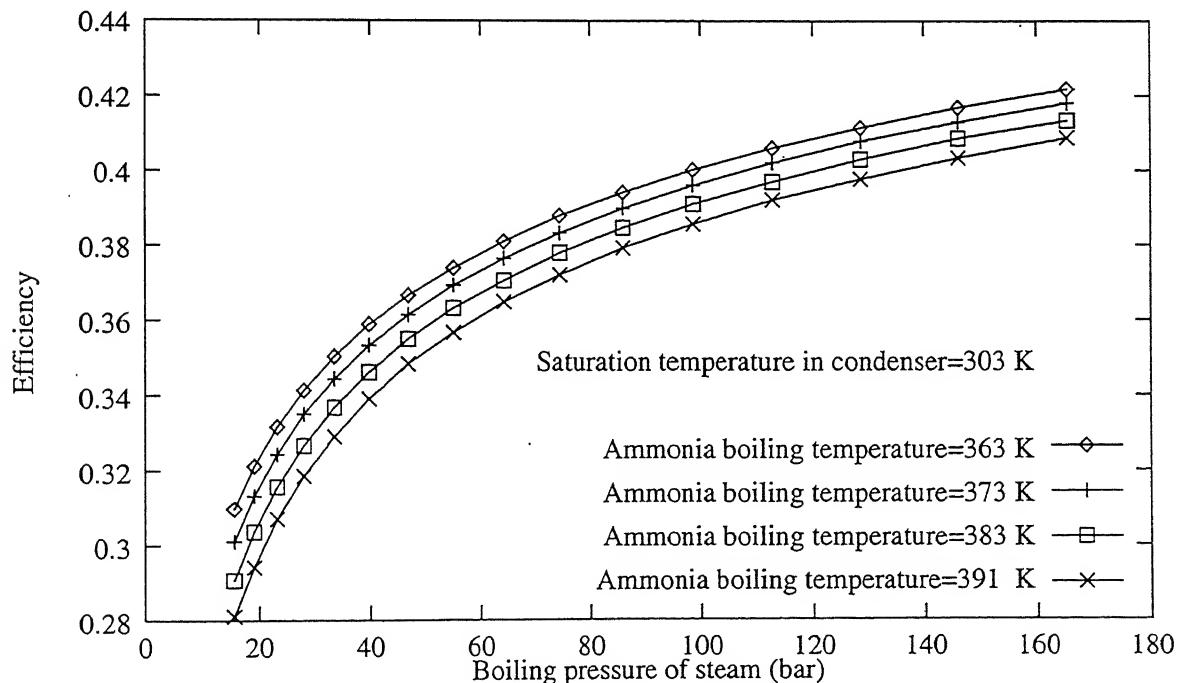


Figure 4.9: Variation in efficiency of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

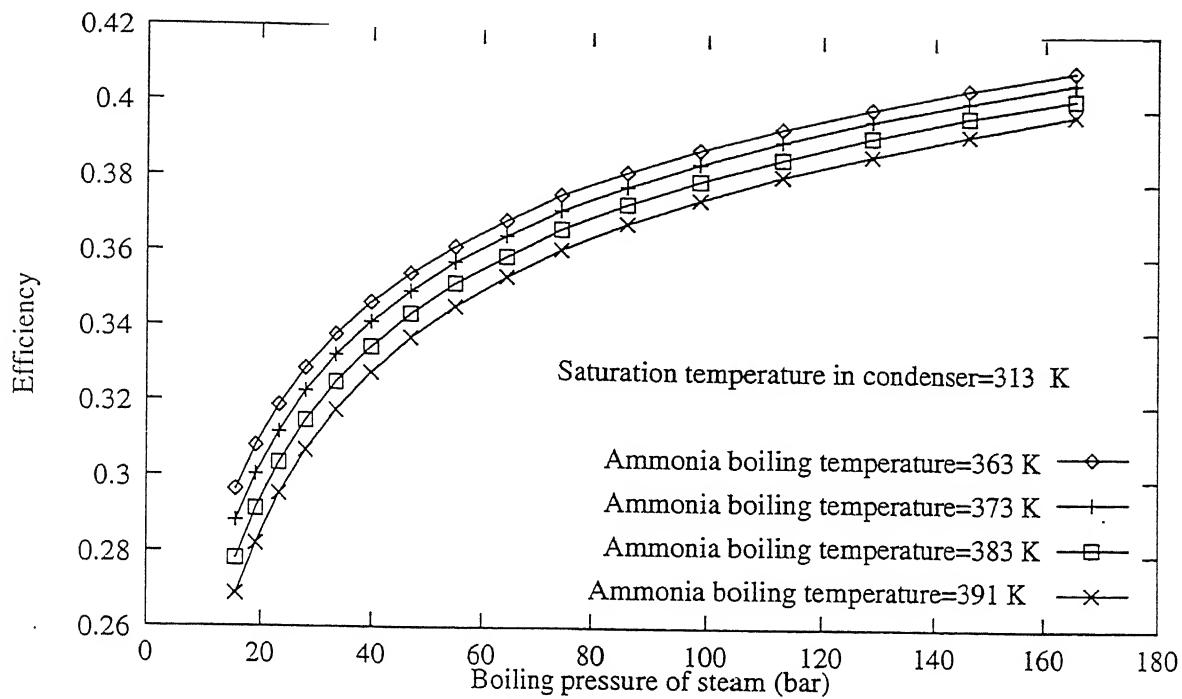


Figure 4.10: Variation in efficiency of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = 313 K

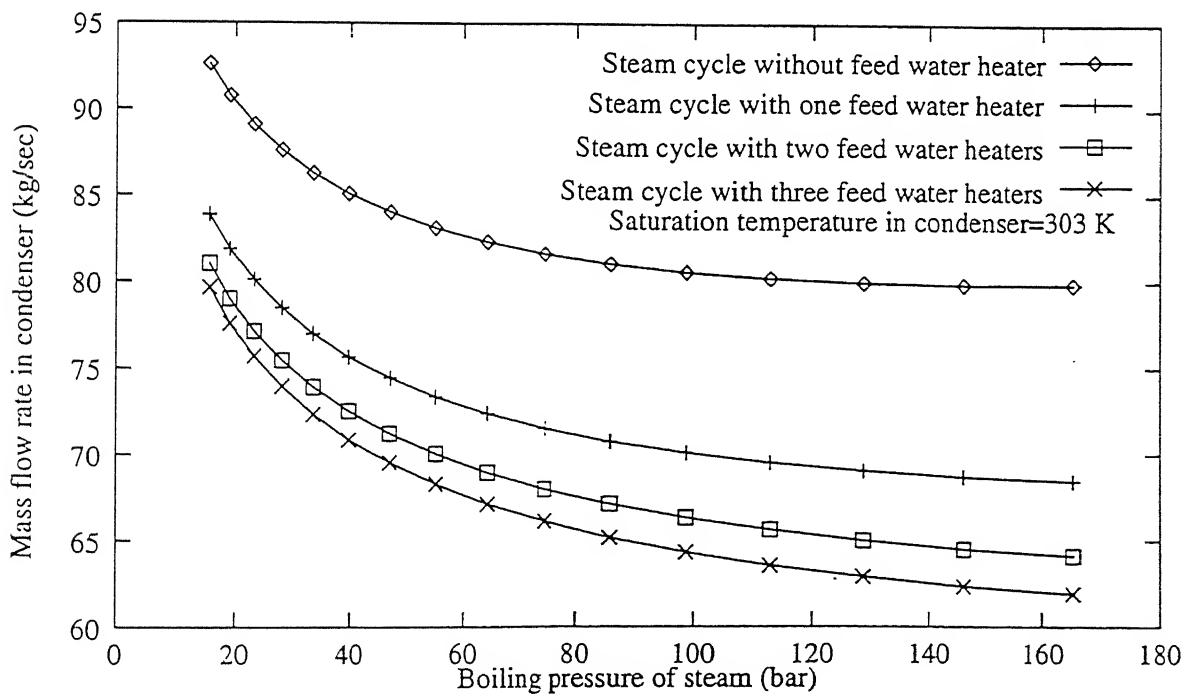


Figure 4.11: Variation in mass flow rate in condenser of steam cycle with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

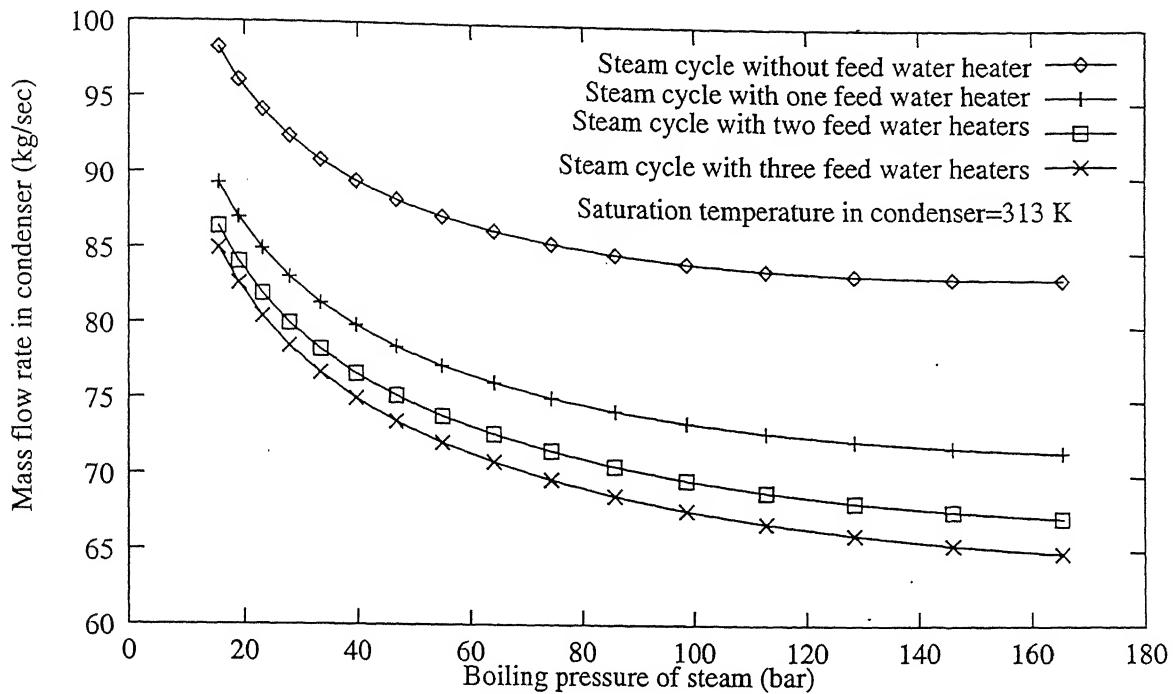


Figure 4.12: Variation in mass flow rate in condenser with boiling pressure of steam at the constant condenser temperature = 313 K

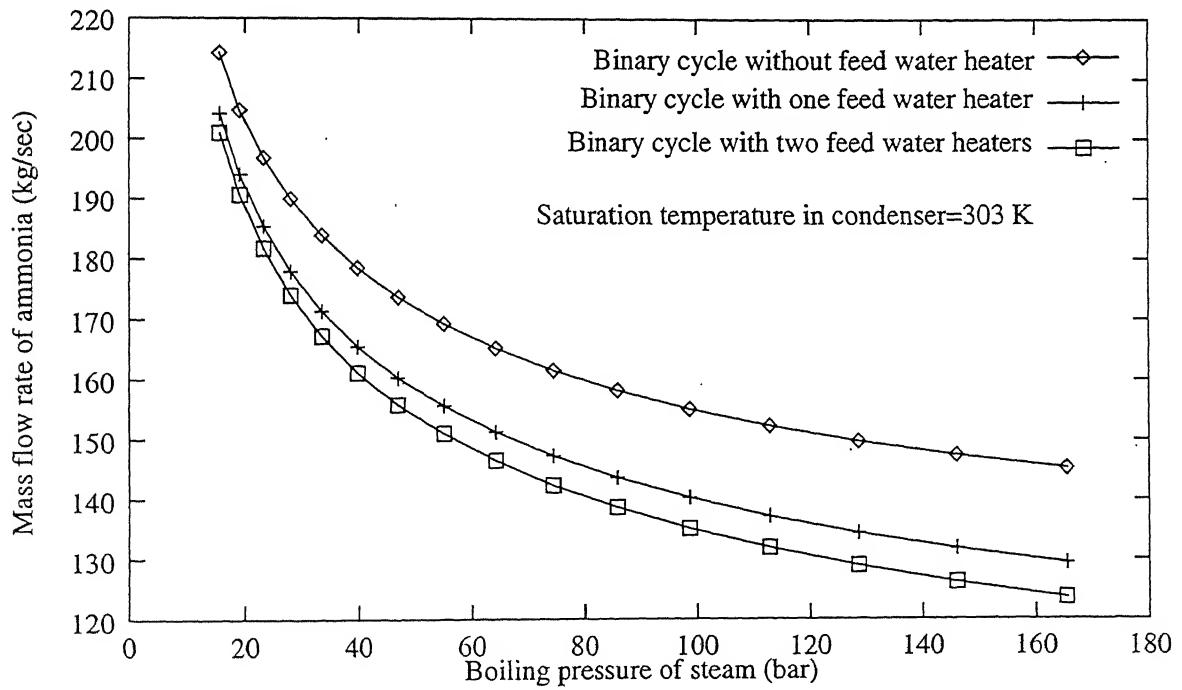


Figure 4.13: Variation in mass flow rate of ammonia of binary-vapour cycle with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

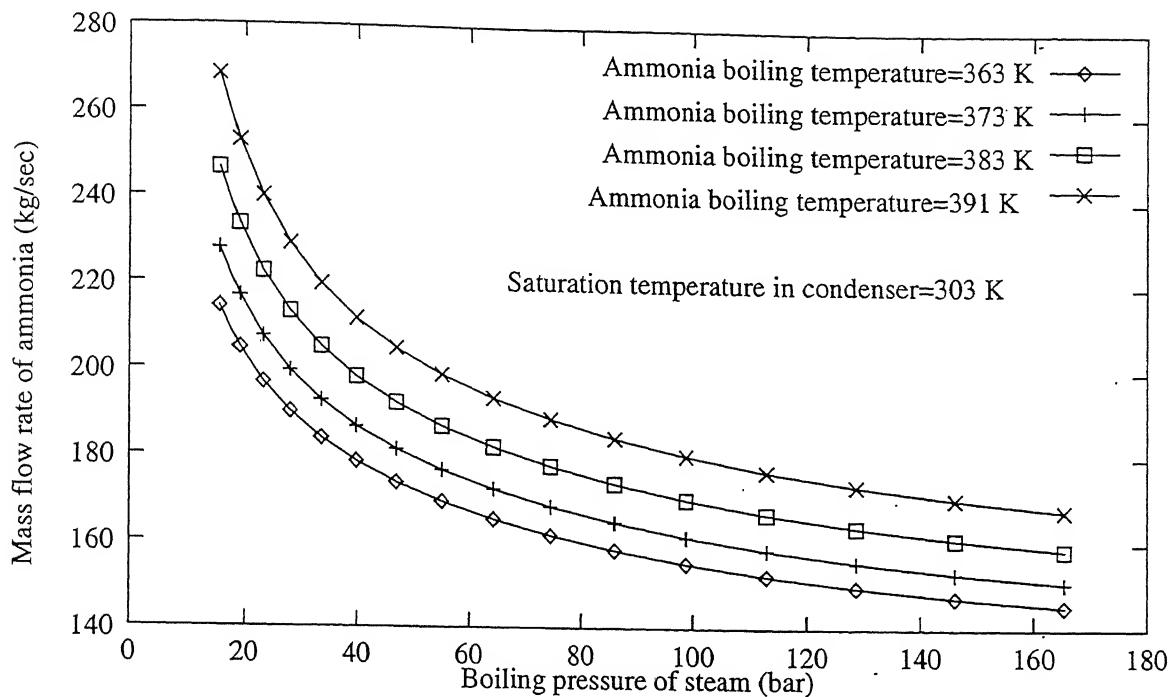


Figure 4.14: Variation in mass flow rate of ammonia of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

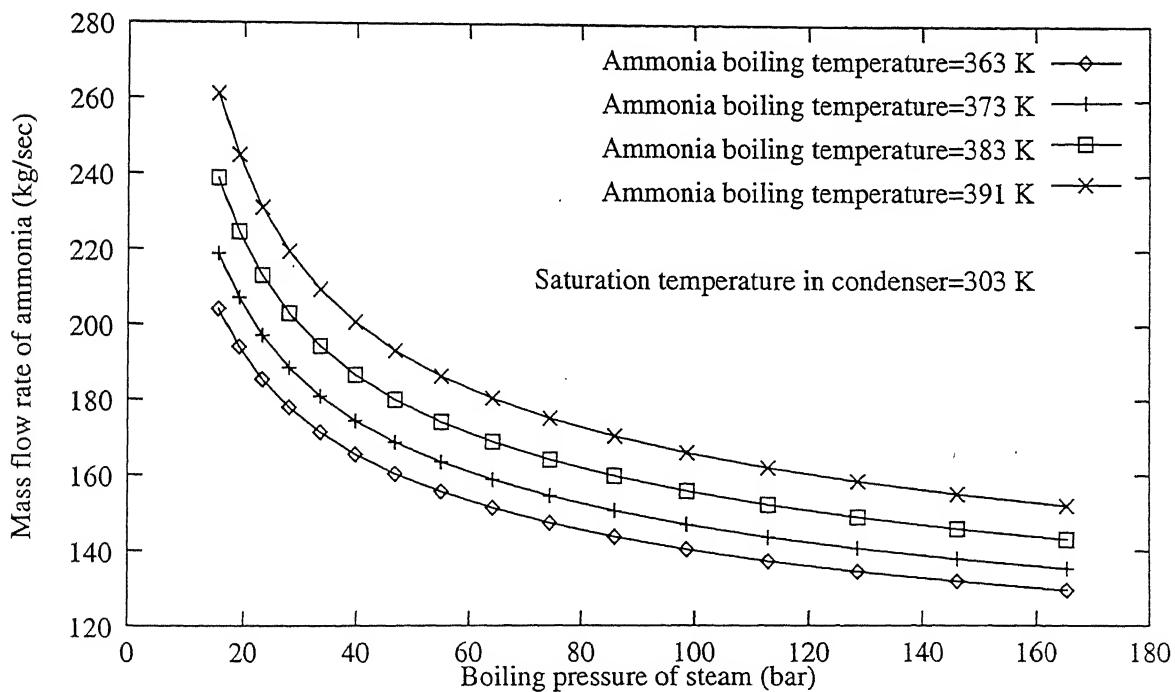


Figure 4.15: Variation in mass flow rate of ammonia of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

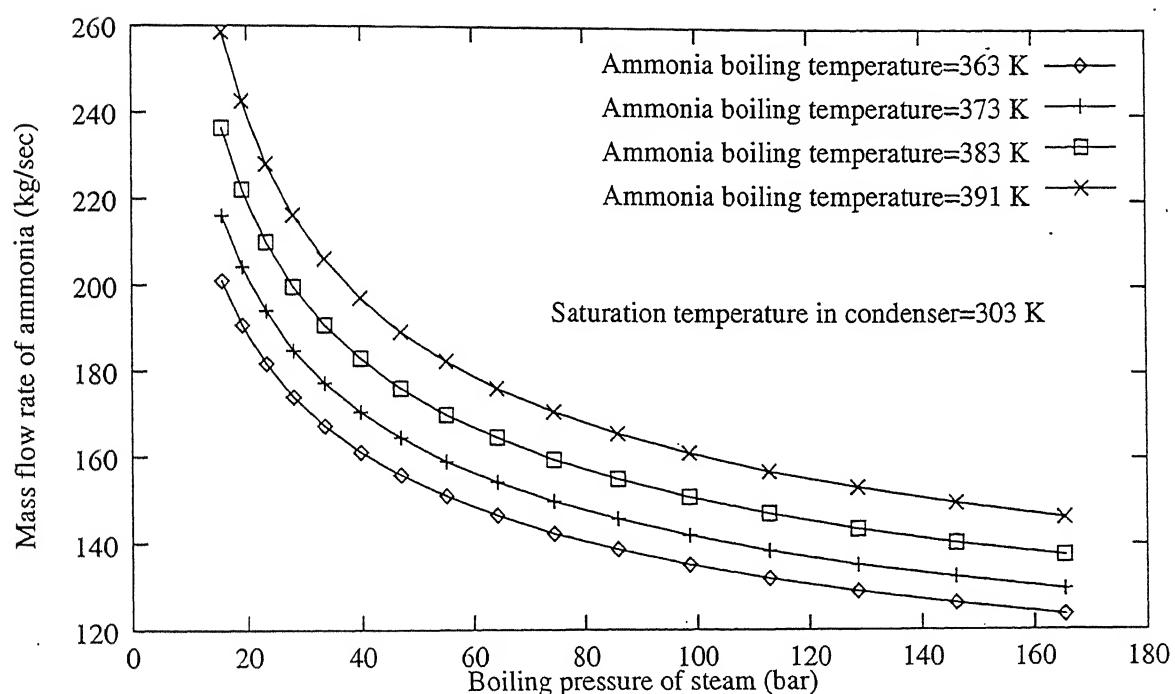


Figure 4.16: Variation in mass flow rate of ammonia of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

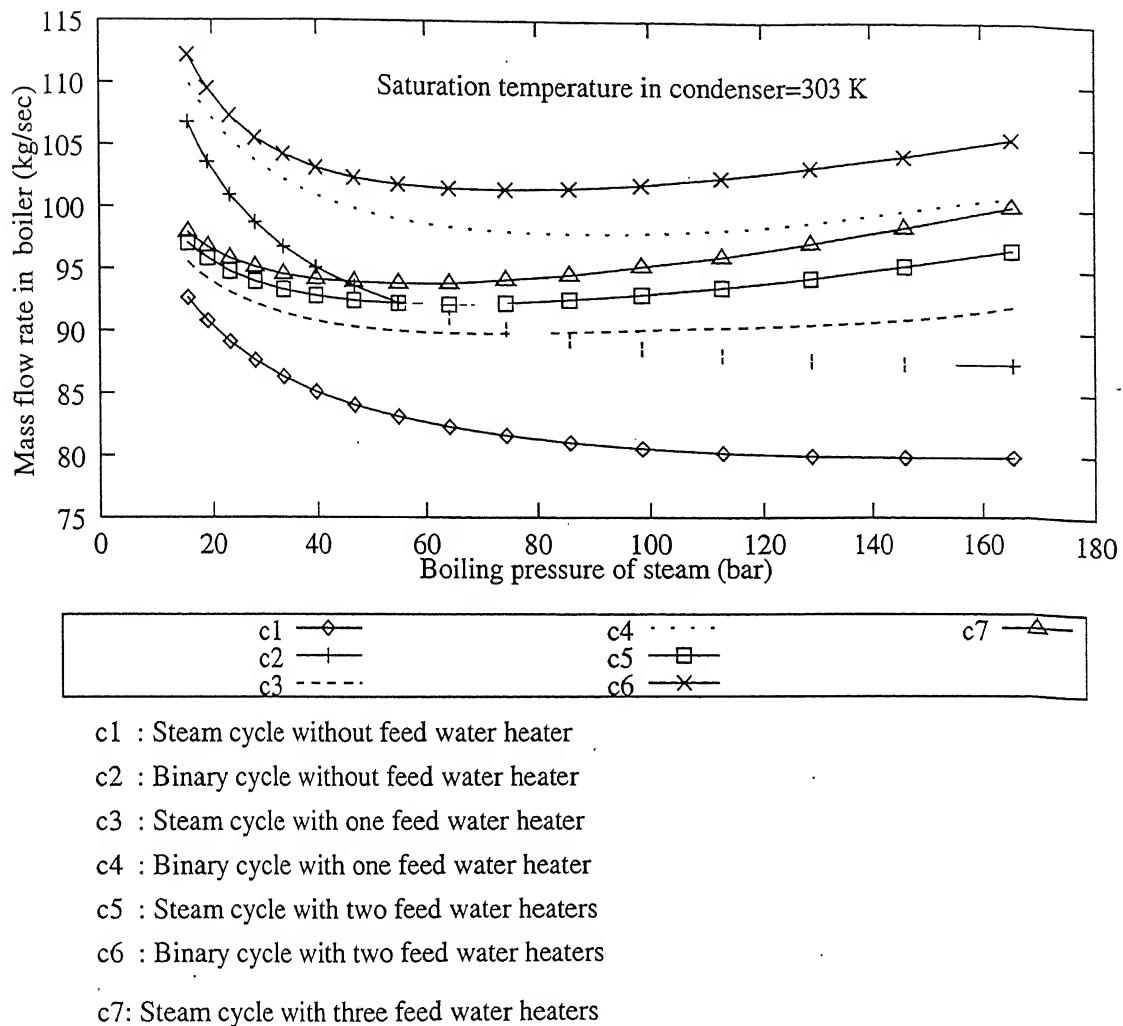


Figure 4.17: Variation in mass flow rate in boiler with boiling pressure of steam at the constant saturation temperature in condenser =303 K

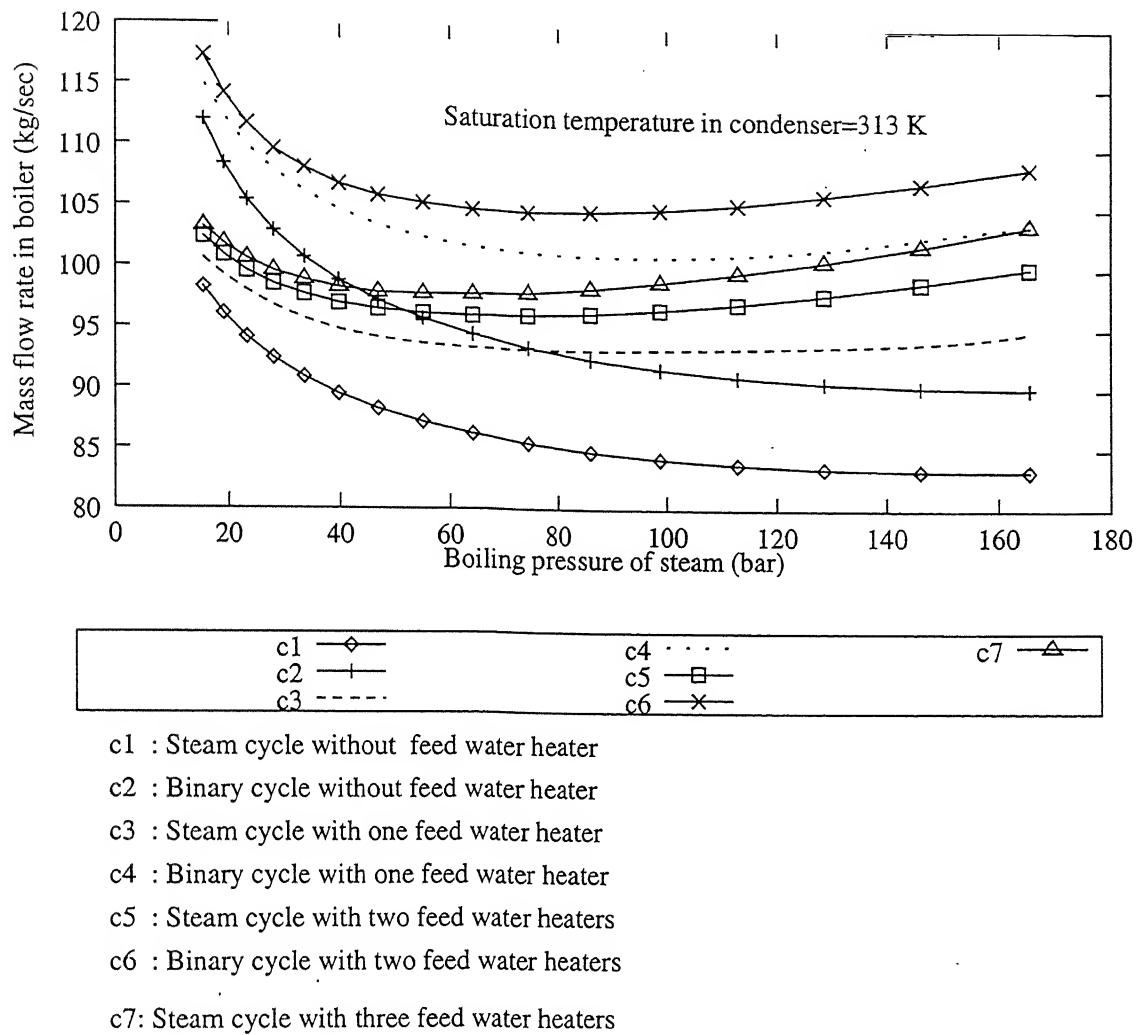


Figure 4.18: Variation in mass flow rate in boiler with boiling pressure of steam at the constant saturation temperature in condenser = 313 K

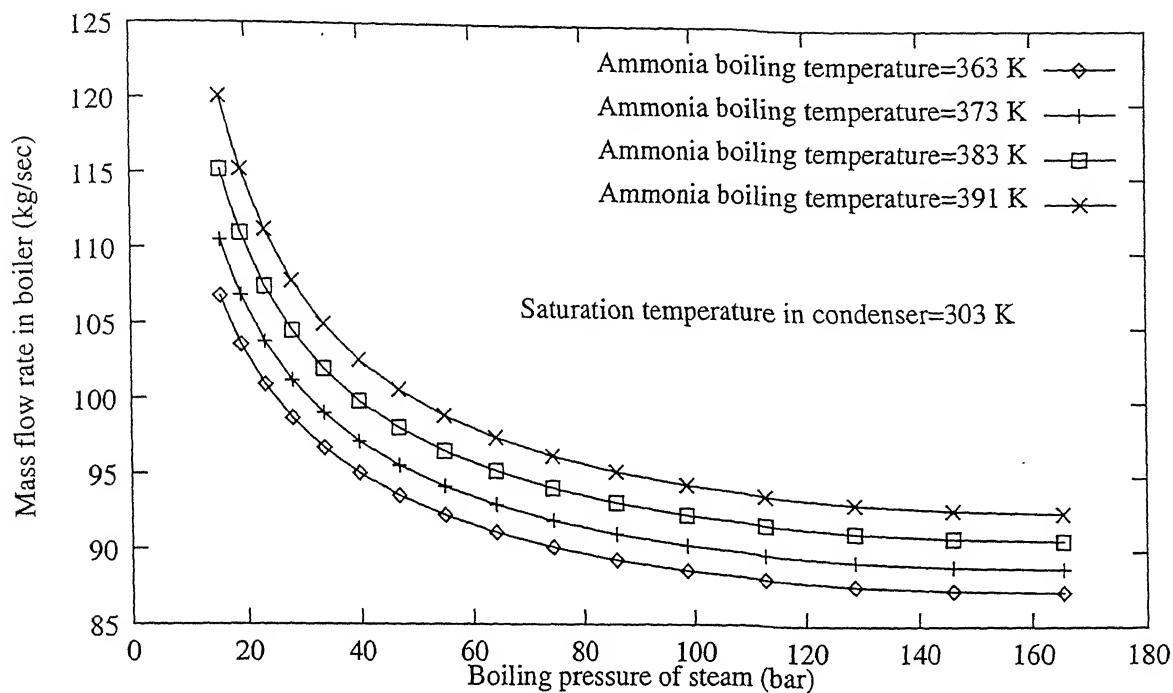


Figure 4.19: Variation in mass flow rate in boiler of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

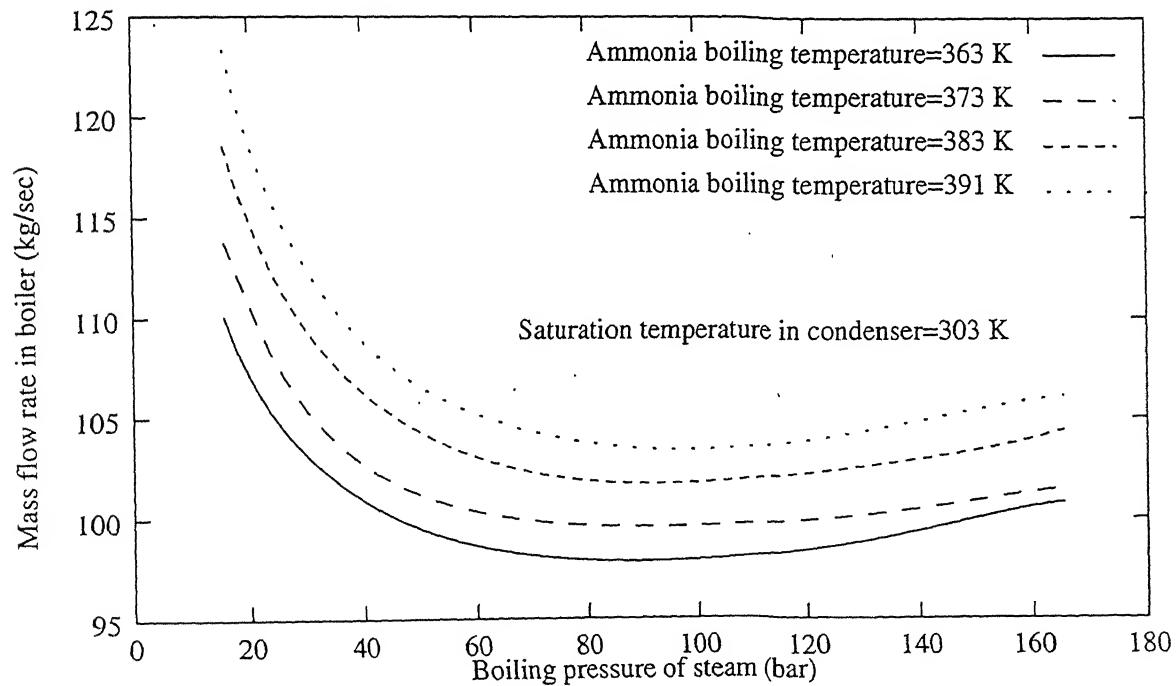


Figure 4.20: Variation in mass flow rate in boiler of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

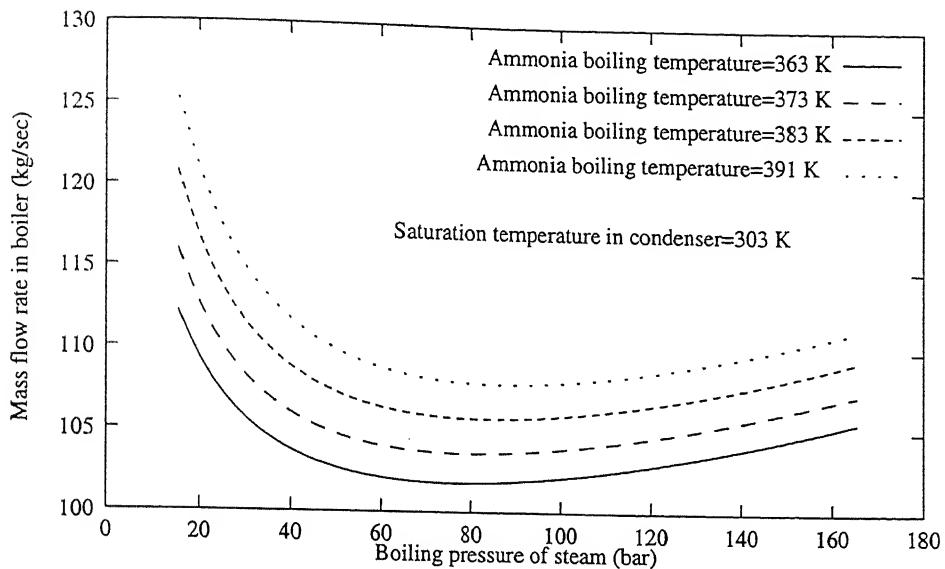


Figure 4.21: Variation in mass flow rate in boiler of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

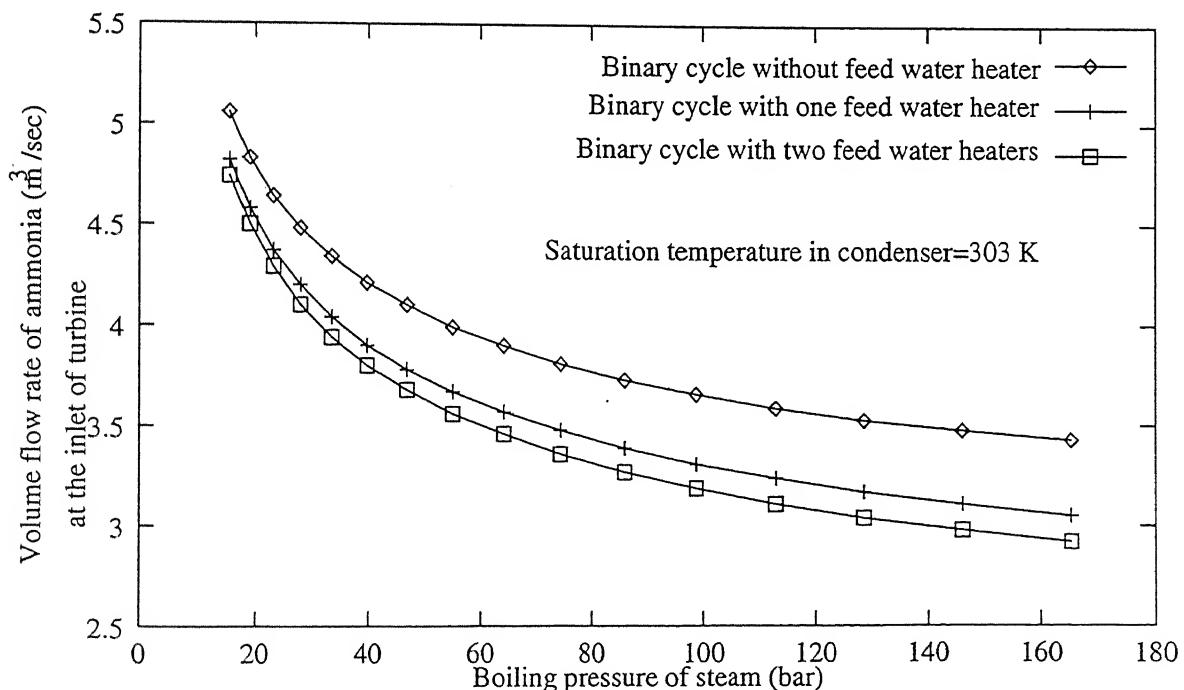


Figure 4.22: Variation in volume flow rate of ammonia with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

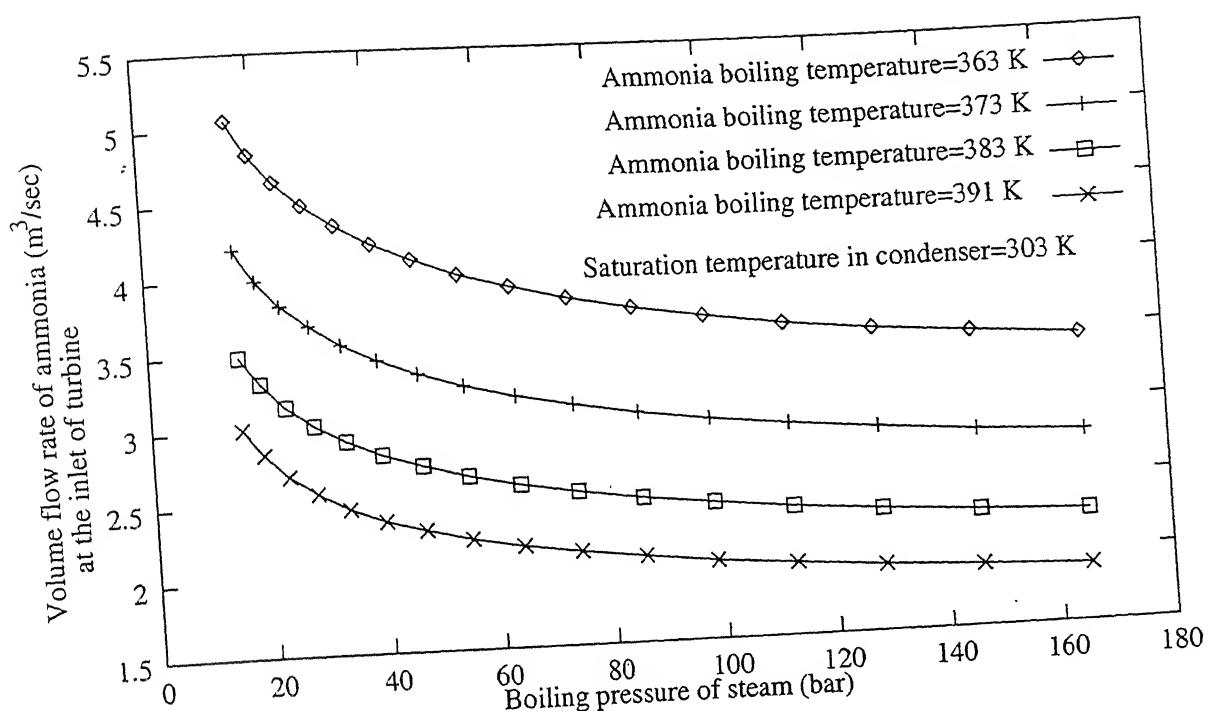


Figure 4.23: Variation in volume flow rate of ammonia of binary-vapour cycle without feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

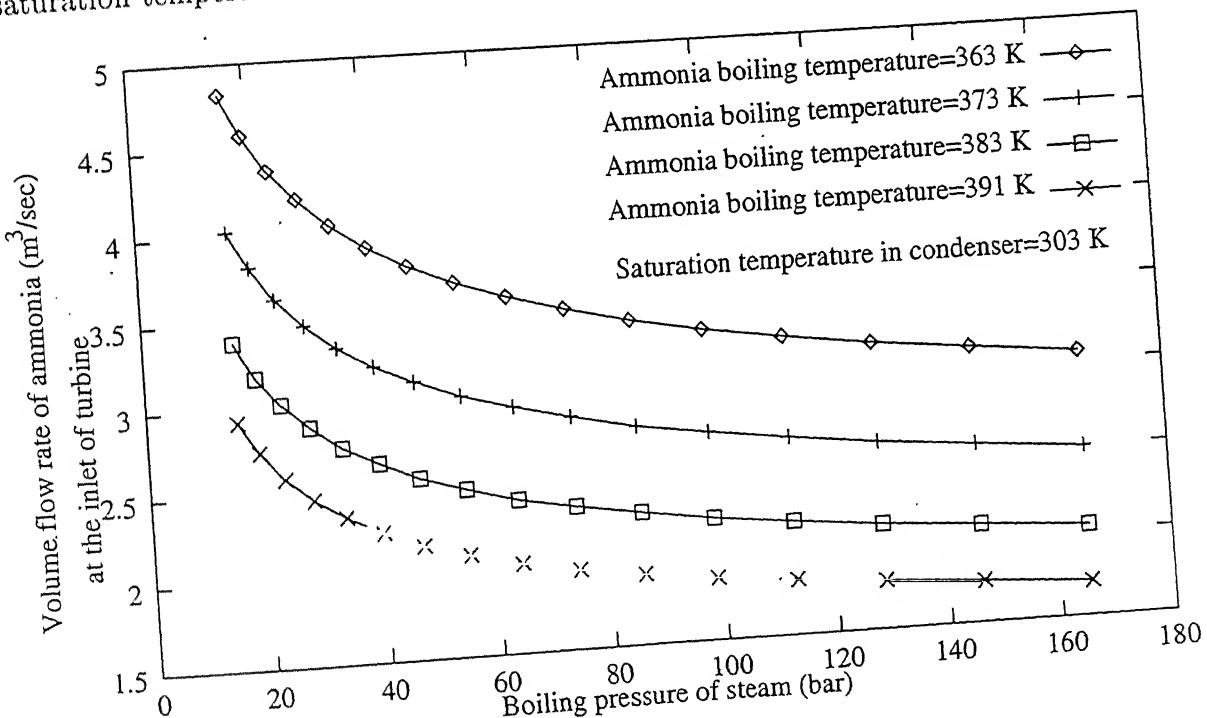


Figure 4.24: Variation in volume flow rate of ammonia of binary-vapour cycle with one feed water heater with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

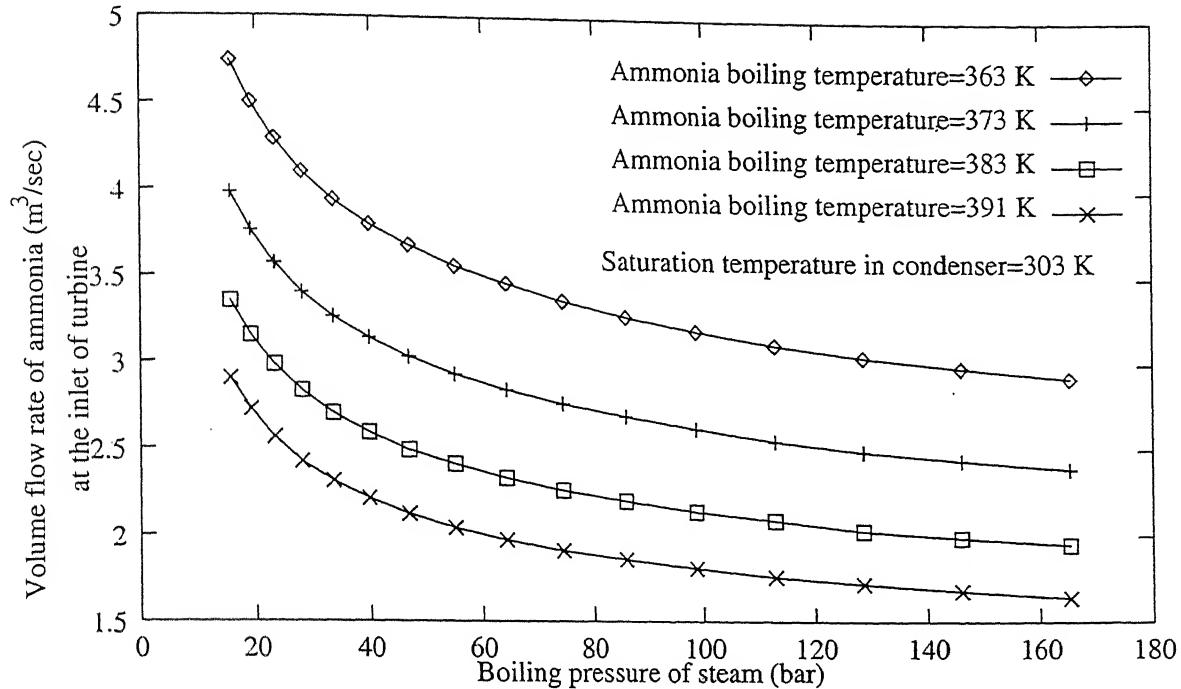


Figure 4.25: Variation in volume flow rate of ammonia of binary-vapour cycle with two feed water heaters with boiling pressure of steam at the constant saturation temperature in condenser = 303 K

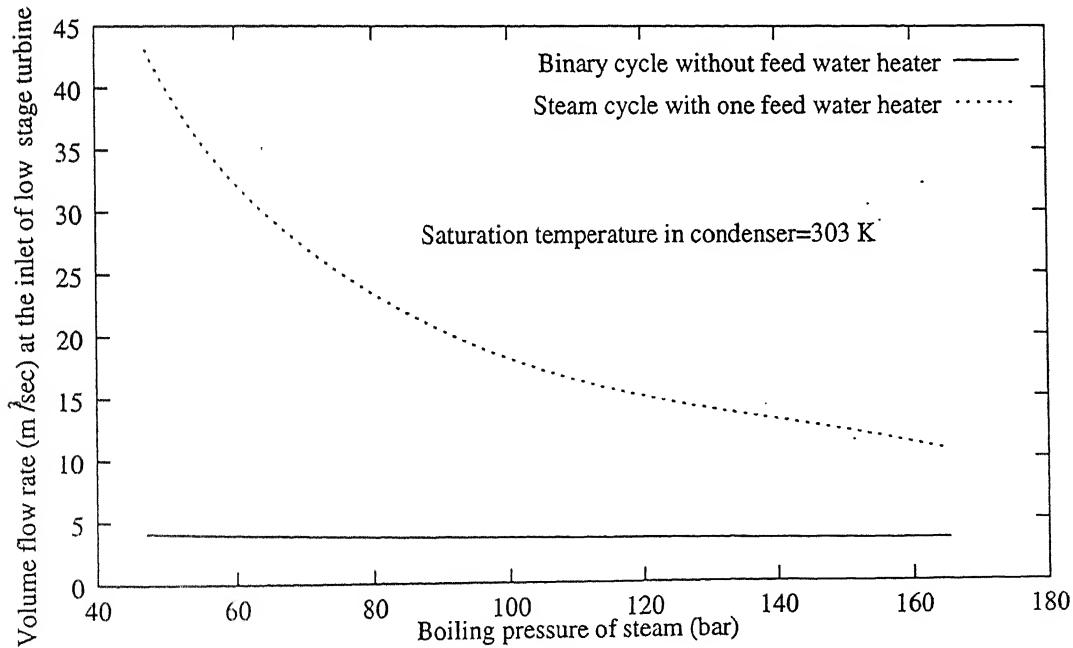


Figure 4.26: Variation in volume flow rate at the inlet of low stage turbine with boiling pressure of the steam at the constant saturation temperature in condenser = 303 K

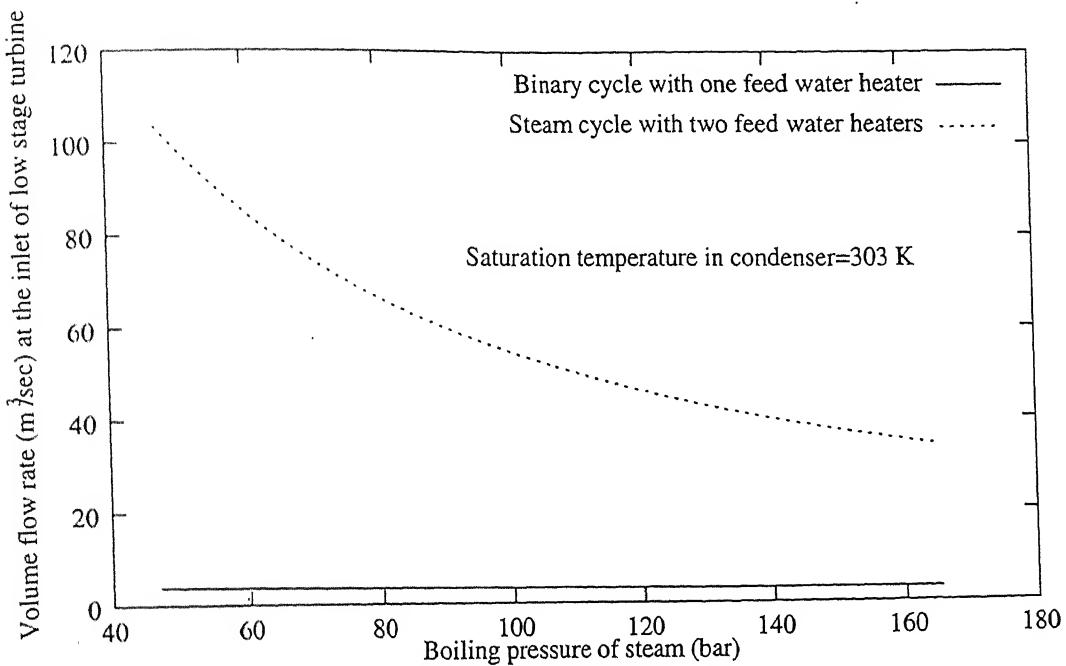


Figure 4.27: Variation in volume flow rate at the inlet of low stage turbine with boiling pressure of the steam at the constant saturation temperature in condenser=303 K

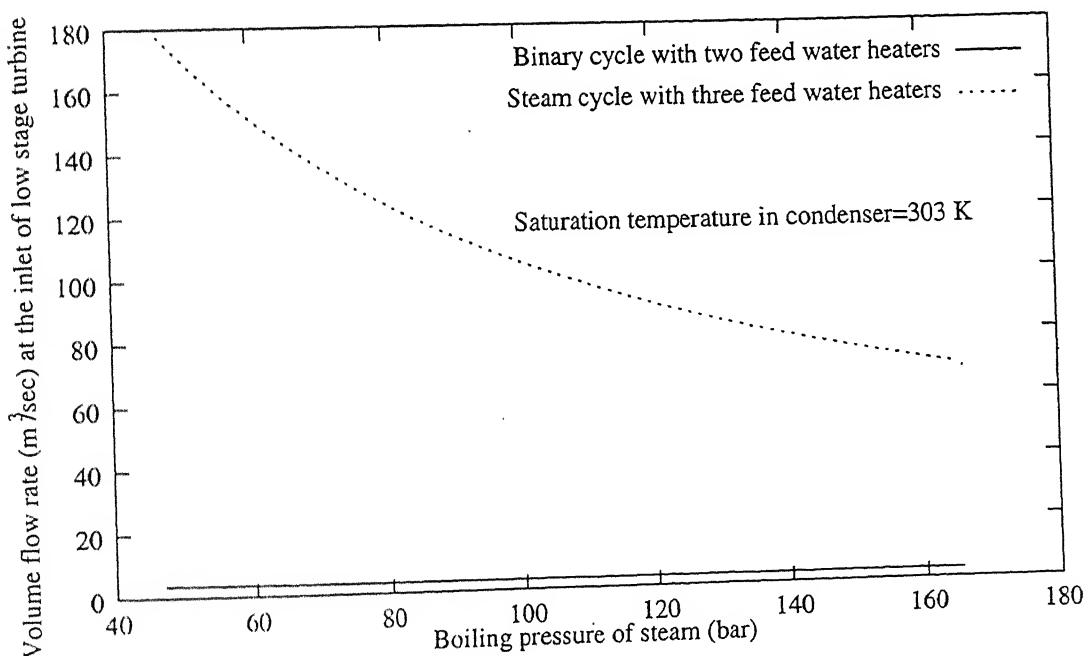


Figure 4.28: Variation in volume flow rate at the inlet of low stage turbine with boiling pressure of the steam at the constant saturation temperature in condenser=303 K

# Chapter 5

## Conclusions and Suggestions

### 5.1 Conclusions

The following conclusions are drawn from the present work :

1. Calculations have been carried out for a 100 MW power plant. In the topping cycle upto 3 feed waters heaters have been considered and the bottoming cycle has no feed heating.
2. Maximum efficiency for a steam cycle with one feed water heater occurs at the arithmetic mean of saturation tempeatures in boiler and condenser.
3. Efficiency of binary-vapour system is slightly less compared to steam cycle. Efficiency reduces as significant temperature difference between steam and ammonia is allowed.
4. Mass flow rate in condenser is much less in ammonia-water binary-vapour cycle than cycle using steam only.
5. Mass flow rate of ammonia decreases as number of feed water heaters are increased. Mass flow rate of ammonia decreases as boiling pressure of

steam increases.

6. At 165.4 bar boiler pressure, volume flow rate at the inlet and exit of the low-pressure turbine is 1/25 and 1/150, respectively in ammonia-water binary-vapour cycle than cycle using steam only.
7. At 165.4 bar boiler pressure, diameters at the inlet and exit of the low-pressure turbine is 1/5 and 1/12, respectively in ammonia-water binary-vapour cycle than cycle using steam only.

## 5.2 Suggestions

1. Effect of feed water heating in the ammonia side needs to be studied in order to see improvement in thermal efficiency of the binary system.
2. Economic analysis can be done for the ammonia-binary-vapour system to find out the initial investment, running cost though it can be intuitively concluded that same will be far less than the cost of system using steam only.
3. Pressure of feed water heaters for power plants running on ammonia binary-vapour system can be selected by using this computer programme.

## Appendix A

### Chebyshev Polynomials

Chebyshev polynomials are used for least square interpolation. Chebyshev polynomials are given by [5].

$$t_0 = 0.5 \quad (\text{A.1})$$

$$t_1 = x \quad (\text{A.2})$$

$$t_2 = 2x^2 - 1 \quad (\text{A.3})$$

$$t_{n+1} = 2t_1 t_n - t_{n-1} \quad (\text{A.4})$$

## Appendix B

# Thermodynamic Properties of Ammonia

### B.1 Specific Volume of Saturated Liquid of Ammonia

Specific volume of saturated liquid of ammonia is given by [11].

$$v_f = \frac{1}{q_l} \quad (B.1)$$

$$q_l = q_c * \left[ 1 + \sum_{i=1}^5 c_i \left( 1 - \frac{T}{T_c} \right)^{\frac{i}{3}} \right] \quad (B.2)$$

Where,  $v_f$  = specific volume of liquid ( $m^3/kg$ )

$q_c$  = critical liquid density =  $234.0\ kg/m^3$

$q_l$  = density of saturated liquid ( $kg/m^3$ )

$c_1 = 1.615217$

$c_2 = 1.807593$

$c_3 = -1.401165$

$c_4 = 0.9919760$

$$c_5 = -0.03569048$$

$T$ =saturation temperature (K)

$T_c$  =critical temperature = 405.50 K

## B.2 Specific Volume of Saturated Vapour of Ammonia

$$v_g = \sum_{n=0}^7 a_n * t_n(x) * 10^{-3} \quad - \quad (B.3)$$

Where,  $v_g$  =specific volume of saturated vapour ( $m^3/kg$ )

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature ( $^{\circ}C$ )

	$0 \leq t_s < 60$	$60 \leq t_s \leq 120$
$x$	$\frac{t_s - 30}{30}$	$\frac{t_s - 90}{30}$
$a_0$	277.6908	53.3893
$a_1$	-113.6503	-18.8436
$a_2$	29.1961	3.1518
$a_3$	-6.02	-0.5403
$a_4$	1.0727	0.0509
$a_5$	-0.1726	-0.0118
$a_6$	0.0254	-0.0004
$a_7$	-0.0035	-0.0006

## B.3 Entropy of Saturated Liquid of Ammonia

$$s_f = \sum_{n=0}^7 a_n * t_n(x) \quad - \quad (B.4)$$

Where,  $s_f$  =entropy of saturated liquid of ammonia (kJ/kg K)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature (°C)

	$0 \leq t_s < 60$	$60 \leq t_s \leq 120$
x	$\frac{t_s - 30}{30}$	$\frac{t_s - 90}{30}$
$a_0$	4.7461	6.6062
$a_1$	0.4632	0.4911
$a_2$	-0.00724	0.0243
$a_3$	0.0008	0.0079
$a_4$	0.0001	0.002
$a_5$	0.0	0.0006
$a_6$	0.0	0.0002
$a_7$	0.0	0.0001

#### B.4 Entropy of Saturated Vapour of Ammonia

$$s_g = \sum_{n=0}^7 a_n * t_n(x) \quad (B.5)$$

Where,  $s_g$  =entropy of saturated vapour of ammonia (kJ/kg K)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature (°C)

	$0 \leq t_s < 60$	$60 \leq t_s \leq 120$
$x$	$\frac{t_s - 30}{30}$	$\frac{t_s - 90}{30}$
$a_0$	12.3488	11.0461
$a_1$	-0.3308	-0.3505
$a_2$	0.0118	-0.0272
$a_3$	-0.0016	-0.0092
$a_4$	0.0	-0.0024
$a_5$	0.0	-0.0007
$a_6$	0.0	-0.0002
$a_7$	0.0	-0.0001

## B.5 Enthalpy of Saturated Liquid of Ammonia

$$h_f = \sum_{n=0}^7 a_n * t_n(x) \quad (B.6)$$

	$0 \leq t_s < 60$	$60 \leq t_s \leq 120$
$x$	$\frac{t_s - 30}{30}$	$\frac{t_s - 90}{30}$
$a_0$	1128.7042	1766.23
$a_1$	142.0403	185.9883
$a_2$	1.6345	13.7632
$a_3$	0.2154	3.2727
$a_4$	0.0299	0.85
$a_5$	0.0019	0.2493
$a_6$	-0.0016	0.0756
$a_7$	-0.0026	0.0247

Where,  $h_f$  = enthalpy of saturated liquid of ammonia (kJ/kg)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature (°C)

## B.6 Enthalpy of Saturated Vapour of Ammonia

$$h_g = \sum_{n=0}^7 a_n * t_n(x) \quad (B.7)$$

Where,  $h_g$  =enthalpy of saturated vapour of ammonia (kJ/kg )

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature (°C)

	$0 \leq t_s < 54$	$54 \leq t_s \leq 120$
$x$	$\frac{t_s - 27}{27}$	$\frac{t_s - 87}{33}$
$a_0$	3408.1289	3358.1124
$a_1$	15.8344	-53.4833
$a_2$	-3.5502	-19.7206
$a_3$	-0.1692	-4.2896
$a_4$	-0.0138	-1.2255
$a_5$	-0.0053	-0.3813
$a_6$	0.0	-0.1242
$a_7$	0.0	-0.0406

## B.7 Saturation Pressure of Ammonia

Saturation pressure of ammonia is given by [11].

$$\ln\left(\frac{p}{p_c}\right) = A + B\left(\frac{T}{T_c}\right) + \frac{C}{\frac{T}{T_c}} + D\left(\frac{T}{T_c}\right)^2 + E\left(\frac{T}{T_c}\right)^3 + \frac{F\left(1 - \frac{T}{T_c}\right)^{1.5}}{\left(\frac{T}{T_c}\right)} \quad (B.8)$$

Where,  $p$ =saturation pressure (bar)

$$A=19.667984$$

$$B=-15.549930$$

$$C=-11.079219$$

$$D=9.114107$$

$$F=1.81269$$

$$T_c = \text{critical temperature} = 405.50 \text{ K}$$

$$p_c = \text{critical pressure} = 113.53 \text{ bar}$$

$$T = \text{saturation temperature ( K )}$$

### B.8 Enthalpy of Superheated Vapour of Ammonia

$$h_g = \sum_{n=0}^7 c_n(y) * t_n(x) \quad (B.9)$$

$$y = \frac{t_s - 60}{60} \quad (B.10)$$

$$x = \frac{degup - 90}{90} \quad (B.11)$$

$$c_n(y) = \sum_{m=0}^7 b_m * t_m(y) \quad (B.12)$$

Where,  $h_g$  = enthalpy of superheated vapour of ammonia (kJ/kg)

$t_s$  = saturation temperature ( °C)

$degup$  = degree of superheat ( °C)

$t_n(x)$  = Chebyshev polynomials

$t_m(y)$  = Chebyshev polynomials

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	$c_0$	$c_1$	$c_2$	$c_3$	$c_4$	$c_5$	$c_6$	$c_7$
$b_0$	7891.917	533.067	-28.3396	121548.7576	-44241.7932	18589.0214	-7980.04	5876.7
$b_1$	118.4331	62.6531	-16.6766	64371.2718	-28430.3465	13696.8698	-6405.1937	5056.64
$b_2$	-27.7705	11.4393	-5.0403	25641.5942	-14476.6787	8140.9864	-4169.38	3651.23
$b_3$	-3.9125	2.4247	-1.7067	10563.6417	-6933.6277	4324.8335	-2381.36	2301.39
$b_4$	-0.9412	0.7578	-0.0692	4239.3495	-3138.5855	2143.2862	-1260.83	1334.27
$b_5$	-0.311	0.2387	-0.2306	1664.5164	-1371.9216	1008.3361	-623.62	720.25
$b_6$	-0.1065	0.0801	-0.0875	653.8141	-578.16555	445.869	-294.19	361.09
$b_7$	-0.0427	0.0307	-0.0365	272.0216	-252.5554	209.4787	-139.78	187.36

### B.9 Entropy of Superheated Vapour of Ammonia

$$s_g = \sum_{n=0}^7 a_n(y) * t_n(x) * 10^{-5} \quad (B.13)$$

$$x = \frac{degup - 90}{90} \quad (B.14)$$

$$a_n(y) = \sum_{m=0}^7 b_m * t_m(y) \quad (B.15)$$

Where,  $s_g$  = entropy of superheated vapour of ammonia (kJ/kg K)

$t_s$  = saturation temperature ( °C)

$degup$  = degree of superheat ( °C)

$t_n(x)$  = Chebyshev polynomials

$t_m(y)$  = Chebyshev polynomials

	$a_0$	$a_1$	$a_2$	$a_3$	$a_4$
$t_s$	$0 \leq t_s \leq 120$	$18 \leq t_s \leq 120$	$0 \leq t_s \leq 120$	$0 \leq t_s \leq 120$	$0 \leq t_s \leq 66$
$y$	$\frac{t_s-60}{60}$	$\frac{t_s-69}{51}$	$\frac{t_s-60}{60}$	$\frac{t_s-60}{60}$	$\frac{t_s-33}{33}$
$b_0$	2634039.2858	128871.3092	-13627.0575	3795.1983	-547.0675
$b_1$	-107112.4378	6422.3971	-3485.8579	1553.5033	-120.9569
$b_2$	7493.0357	2014.5929	-1149.7532	609.5956	-25.3876
$b_3$	-2174.082	392.8705	-379.3708	248.7585	-96.813
$b_4$	-82.0241	126.0958	-147.5753	96.2697	-5.7800
$b_5$	-86.563	37.4694	-54.0439	37.2890	-4.0002
$b_6$	-25.6332	12.4024	-18.5529	18.6148	-2.4960
$b_7$	-11.1804	4.7725	-7.3680	8.7564	-1.4394

	$a_4$	$a_5$	$a_5$	$a_6$	$a_7$
$t_s$	$70 \leq t_s \leq 120$	$0 \leq t_s \leq 66$	$70 \leq t_s \leq 120$	$70 \leq t_s \leq 120$	$70 \leq t_s \leq 120$
$y$	$\frac{t_s-95}{25}$	$\frac{t_s-33}{33}$	$\frac{t_s-95}{25}$	$\frac{t_s-95}{25}$	$\frac{t_s-95}{25}$
$b_0$	-1964.0550	1467.8056	827.2694	-354.0112	255.0752
$b_1$	-726.4124	39.3903	394.6464	-198.2926	168.7846
$b_2$	-233.8985	10.0221	148.9912	-83.2312	82.2962
$b_3$	-69.6397	3.7600	49.6649	-30.0976	33.7420
$b_4$	-20.2421	2.1239	15.6666	-9.9786	12.5981
$b_5$	-6.0054	1.3134	4.8553	-3.1896	4.3869
$b_6$	-1.6103	0.8349	1.4854	-0.9729	1.4739
$b_7$	-0.5922	0.5966	0.4882	-0.1225	0.5743

$t_s$	$a_1$
0	59824.0811
6	59731.4234
12	59701.0446
18	59731.7476

$t_s$	$a_4$	$a_5$
66	-443.3623	131.4483
70	-435.1302	130.1301

$t_s$	$a_6$	$a_7$
0	-9.7159	2.884
6	-10.6687	3.078
12	-11.6208	3.3712
18	-12.7873	4.0379
24	-15.0987	4.0384
30	-15.8894	5.0300
36	-17.6511	5.5724
42	-19.7128	6.7977
48	-22.1636	7.8696
54	-25.7231	9.4115
60	-29.8201	11.1494
66	-20.5612	68.1460
70	-39.4898	16.4504

## B.10 Specific Volume of Superheated Vapour of Ammonia

$$v_g = \sum_{n=0}^7 p_n(y) * t_n(x) * 10^{-3} \quad (B.16)$$

$$x = \frac{deg_{sup} - 90}{90} \quad (B.17)$$

$$p_n(y) = \sum_{m=0}^7 b_m * t_m(y) \quad (B.18)$$

Where,  $v_g$  = specific volume of superheated *vapour* of ammonia ( $m^3/kg$ )

$t_s$  = saturation temperature ( °C)

$deg_{sup}$  = degree of superheat

$t_n(x)$  = Chebyshev polynomials

$t_m(y)$  = Chebyshev polynomials

	$p_0$	$p_1$	$p_2(10^{-4})$	$p_2(10^{-4})$
$t_s$	$0 \leq t_s \leq 120$	$0 \leq t_s \leq 120$	$0 \leq t_s \leq 95$	$100 \leq t_s \leq 120$
$y$	$\frac{t_s - 60}{60}$	$\frac{t_s - 60}{60}$	$\frac{t_s - 47.5}{47.5}$	$\frac{t_s - 110}{10}$
$b_0$	537.829	76.3138	-28833.858	-20798.8
$b_1$	-329.0478	-43.6345	5949.7742	-466.3333
$b_2$	138.2894	18.8322	-1733.7865	-118.4286
$b_3$	-48.0197	-6.6411	220.1486	-11.6667
$b_4$	14.6213	2.0878	-50.0565	0.0
$b_5$	-4.0635	-0.5904	3.117	0.0
$b_6$	1.0371	0.1558	-1.575	0.0
$b_7$	-0.2494	-0.0375	-0.827	0.0

	$p_3(10^{-4})$	$p_3(10^{-4})$	$p_4(10^{-4})$	$p_4(10^{-4})$
$t_s$	$0 \leq t_s \leq 80$	$80 \leq t_s \leq 120$	$0 \leq t_s \leq 60$	$75 \leq t_s \leq 120$
$y$	$\frac{t_s-40}{40}$	$\frac{t_s-100}{20}$	$\frac{t_s-30}{30}$	$\frac{t_s-97.5}{22.5}$
$b_0$	7883.0203	6610.0180	-2127.9431	-2404.0019
$b_1$	-1463.3189	586.8857	319.6638	-485.1068
$b_2$	447.0288	193.5838	-85.2628	-155.0146
$b_3$	-49.6254	33.0	8.1725	-37.1587
$b_4$	11.3265	7.4414	-1.5568	-9.8119
$b_5$	-1.8048	1.3333	-0.1177	-1.5911
$b_6$	0.8015	-0.0356	0.2433	-1.1893
$b_7$	0.2764	-1.219	-0.2208	-0.1431

	$p_5(10^{-4})$	$p_5(10^{-4})$	$p_6(10^{-4})$	$p_7$
$t_s$	$0 \leq t_s \leq 42$	$75 \leq t_s \leq 120$	$75 \leq t_s \leq 120$	$75 \leq t_s \leq 120$
$y$	$\frac{t_s-21}{21}$	$\frac{t_s-97.5}{22.5}$	$\frac{t_s-97.5}{22.5}$	$\frac{t_s-97.5}{22.5}$
$b_0$	559.6989	991.3737	-417.4901	0.0290
$b_1$	-67.1546	313.8334	-170.5668	0.0155
$b_2$	12.7195	104.061	-61.4864	0.0065
$b_3$	-2.4791	28.5851	-17.9669	0.0023
$b_4$	-0.0423	8.5417	-5.0854	0.0007
$b_5$	0.7694	2.3931	-1.6757	0.0002
$b_6$	2.4734	1.2074	-0.1982	0.0
$b_7$	5.8643	0.6985	-0.2964	0.0

$t_s$	$p_2$	$t_s$	$p_4$	$t_s$	$p_5$
95	-1.003	60	-0.0823	42	0.0232
100	-1.004	66	-0.0822	48	0.023
		70	-0.0822	54	0.0232
		75	-0.0844	60	0.0235
				66	0.0244
				70	0.0262
				75	0.0264

$t_s$	$p_6(10^{-4})$	$p_7$
0	-85	0.0024
6	-79	0.0021
12	-70	0.0023
18	-66	0.0020
24	-69	0.0024
30	-62	0.0021
36	-62	0.0020
42	-60	0.0021
48	-64	0.0023
54	-66	0.0022
60	-69	0.0026
66	-72	0.0031
70	-67	0.0040
75	-85	0.0037

# Appendix C

## Thermodynamic Properties of Steam

### C.1 Enthalpy of Saturated Liquid

$$h_f = \sum_{n=0}^7 a_n * t_n(x) * 0.1 \quad (C.1)$$

	$0 \leq t_s < 100$	$100 \leq t_s < 236$	$236 \leq t_s < 320$	$320 \leq t_s \leq 370$
$x$	$\frac{t_s - 50}{50}$	$\frac{t_s - 168}{68}$	$\frac{t_s - 278}{42}$	$\frac{t_s - 345}{25}$
$a_0$	4188.1623	14290.9635	24664.8121	33037.1912
$a_1$	2094.0919	2991.7219	2209.0742	2069.2907
$a_2$	1.31271	42.2673	70.2006	223.172
$a_3$	1.29107	5.929	8.5709	69.4128
$a_4$	-0.23937	0.609	1.2298	26.6501
$a_5$	0.11589	0.0659	0.2662	10.6612
$a_6$	-0.05306	-0.0214	0.0783	4.5517
$a_7$	0.00071	0.0056	0.1157	1.6313

Where,  $h_f$  =enthalpy of saturated liquid (kJ/kg)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature (°C)

## C.2 Enthalpy of Saturated Vapour

$$h_g = \sum_{n=0}^7 a_n * t_n(x) * 0.1 \quad (C.2)$$

Where,  $h_g$  =enthalpy of saturated vapour (kJ/kg)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature (°C)

	$0 \leq t_s < 100$	$100 \leq t_s < 236$	$236 \leq t_s < 320$	$320 \leq t_s \leq 370$
$x$	$\frac{t_s - 50}{50}$	$\frac{t_s - 168}{68}$	$\frac{t_s - 278}{42}$	$\frac{t_s - 345}{25}$
$a_0$	5179.77712	55063.222	55337.0687	51186.9596
$a_1$	87.71343	651.7661	-502.4306	-1712.0818
$a_2$	-1.59094	-135.4466	-151.6567	-395.1305
$a_3$	-0.22595	-13.8817	-16.6366	-113.7785
$a_4$	-0.01904	-0.3957	-2.17	-40.5655
$a_5$	-0.00160	0.2264	-0.1513	-16.1384
$a_6$	0.02274	0.2067	0.2591	-6.2431
$a_7$	0.01425	-0.0237	0.1843	-2.5004

## C.3 Entropy of Saturated Liquid

$$s_f = \sum_{n=0}^7 a_n * t_n(x) * 10^{-3} \quad (C.3)$$

Where,  $s_f$  =entropy of saturated liquid (kJ/kg K)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature ( $^{\circ}\text{C}$ )

	$0 \leq t_s < 100$	$100 \leq t_s < 236$	$236 \leq t_s < 320$	$320 \leq t_s \leq 370$
$x$	$\frac{t_s - 50}{50}$	$\frac{t_s - 168}{68}$	$\frac{t_s - 278}{42}$	$\frac{t_s - 345}{25}$
$a_0$	1357.20862	4008.0589	6106.6591	7491.4305
$a_1$	651.77539	677.0974	390.6981	319.9379
$a_2$	-25.0617	-18.0248	3.7443	31.1037
$a_3$	1.71435	1.9342	1.3434	10.6539
$a_4$	-0.17464	-0.0679	0.0523	4.1465
$a_5$	0.06802	0.0275	-0.0342	1.6835
$a_6$	-0.01755	0.0272	-0.0450	0.5306
$a_7$	-0.00784	-0.0038	0.0031	0.2280

#### C.4 Entropy of Saturated Vapour

$$s_g = \sum_{n=0}^7 a_n * t_n(x) * 10^{-3} \quad (\text{C.4})$$

Where,  $s_g$  =entropy of saturated liquid (kJ/kg K)

$t_n$  =Chebyshev polynomials

$t_s$  =saturation temperature ( $^{\circ}\text{C}$ )

	$0 \leq t_s < 100$	$100 \leq t_s < 236$	$236 \leq t_s < 320$	$320 \leq t_s \leq 370$
$x$	$\frac{t_s - 50}{50}$	$\frac{t_s - 168}{68}$	$\frac{t_s - 278}{42}$	$\frac{t_s - 345}{25}$
$a_0$	16328.7919	13444.9525	11725.5545	10454.1862
$a_1$	-892.0519	-585.0940	-314.9788	-350.0705
$a_2$	89.9277	39.9095	-9.5607	-55.0319
$a_3$	-8.1931	-6.9673	-2.7495	-17.1820
$a_4$	0.7776	0.6076	-0.3603	-6.0641
$a_5$	0.2214	-0.0264	-0.1970	-2.4713
$a_6$	-0.0997	-0.0900	0.1098	-1.0053
$a_7$	0.0214	0.0274	-0.0514	-0.2778

## C.5 Specific Volume of Saturated Liquid

Specific volume of saturated liquid of steam is given by [12].

$$\ln\left(\frac{v_c}{v_f}\right) = a_{11} * \tau^{\frac{1}{3}} + a_{12} * \tau^{\frac{2}{3}} + a_{13} * \tau + a_{14} * \tau^{\frac{7}{3}} + a_{15} * \tau^{\frac{8}{3}} \quad (C.5)$$

$$\tau = \frac{T_c}{T} - 1 \quad (C.6)$$

Where,  $v_c = 0.003147 \text{ (m}^3/\text{kg)}$

$v_f$  = specific volume of saturated liquid ( $\text{m}^3/\text{kg}$ )

$a_{11} = 2.0198239$

$a_{12} = -0.72813257$

$a_{13} = -0.22418367$

$a_{14} = 0.20029413$

$a_{15} = -0.13359233$

$T$  = saturation temperature (K)

$T_c$  = critical temperature = 647.27 K

## C.6 Specific Volume of Saturated Vapour

Specific volume of saturated vapour of steam is given by [12].

$$v_g = v_{fg} + v_f \quad (C.7)$$

$$v_{fg} = \frac{z_{fg} * R * T * 0.001}{p_s} \quad (C.8)$$

$$\ln(1 - z_{fg}) = a41 * \tau^{\frac{1}{3}} + a42 * \tau^{\frac{2}{3}} + a43 * \tau + a44 * \tau^{\frac{5}{3}} + a45 * \tau^3 + a46 * \tau^5 \quad (C.9)$$

$$\tau = \frac{T_c}{T} - 1 \quad (C.10)$$

Where,  $v_g$  = specific volume of saturated liquid ( $m^3/kg$ )

$v_f$  = specific volume of saturated liquid ( $m^3/kg$ )

$R=0.46151\text{kJ/kgK}$

$p_s$  = saturation pressure (MPa)

$a41=-0.80915687$

$a42=-1.4874299$

$a43=-1.8106570$

$a44=-1.4309858$

$a45=0.041862761$

$a46=-0.046394800$

$T$  = saturation temperature (K)

$T_c$  = critical temperature = 647.27 K

## C.7 Saturation Pressure

Saturation pressure of steam is given by [1].

$$\log p = A + B \log z + Cz + Dz \quad (C.11)$$

Where, p=saturation pressure (bar)

$$z=t+273.16=T+0.01$$

$$A=28.59051$$

$$B=-8.2$$

$$C=2.4804 \times 10^{-3}$$

$$D=-3142.31$$

Equation C.11 is valid for the temperatures less than 100°C .

$$\log p = a + \frac{b}{z} + \frac{cx}{z} (10^{dx^2} - 1) + e 10^{fy^{1.25}}$$

$$\log p = a + \frac{b}{z} + \frac{cx}{z} (10^{dx^2} - 1) + e 10^{fy^{1.25}} \quad (C.12)$$

Where, p=saturation pressure (bar)

$$x = z^2 - g$$

$$y = 374.11 - t$$

$$z = t + 273.16 = T + 0.01$$

$$a = 5.432368$$

$$b = -2.0051 \times 10^3$$

$$c = 1.3869 \times 10^{-4}$$

$$d = 1.1965 \times 10^{-11}$$

$$e = -4.4000 \times 10^{-3}$$

$$f = -5.7148 \times 10^{-3}$$

$$g = 2.9370 \times 10^5$$

Equation C.12 is valid for the temperatures above 100°C .

## C.8 Superheated Properties for the Region $0^\circ C \leq t \leq 350^\circ C$

Superheated properties are given by [1].

$$v = \frac{RT}{p} + a_0(\Gamma) + p\Gamma^6 a_1(\Gamma) + (p\Gamma^6)^4 a_2(\Gamma) \quad (C.13)$$

Where,  $v$ =specific volume ( $cm^3/g$ )

$R=4.6152 \text{ bar } cm^3/g \text{ deg K}$

$p$ =saturation pressure (bar)

$T$ =saturation temperature (K)

$$\Gamma = \frac{500}{T} \quad (C.14)$$

$$a_0(\Gamma) = 0.512004 - 1.191807\Gamma + 2.599832\Gamma^2 - 21.433083\Gamma^3 + 15.281761\Gamma^4 - 2.527165\Gamma^5 - 2.454047\Gamma^6 \quad (C.15)$$

$$a_1(\Gamma) = 0.661366 - 3.258346\Gamma + 6.393115\Gamma^2 - 6.447504\Gamma^3 + 3.202128\Gamma^4 - 0.514945\Gamma^5 - 0.120192\Gamma^6 \quad (C.16)$$

$$a_2(\Gamma) * 10^6 = 8.44104 + 28.86344\Gamma - 270.10366\Gamma^2 + 624.08835\Gamma^3 - 675.70455\Gamma^4 + 363.16788\Gamma^5 - 79.26405\Gamma^6 \quad (C.17)$$

$$h - h_0 = \frac{1}{10} \left\{ p \left( a_0 + \Gamma \frac{da_0}{d\Gamma} \right) + \frac{1}{2} p^2 \Gamma^6 \left( 7a_1 + \Gamma \frac{da_1}{d\Gamma} \right) + \frac{1}{5} p^5 \Gamma^{24} \left( 25a_2 + \Gamma \frac{da_2}{d\Gamma} \right) \right\} \quad (C.18)$$

$$h_0 = 1809.25 + 1.48286T + 3.79025 * 10^{-4}T + 46.174 \ln T \quad (C.19)$$

Where, h=enthalpy (KJ/kg)

$$\left( s + \frac{R}{10} \ln p \right) - \left( s + \frac{R}{10} \ln p \right)_0 = \frac{\Gamma}{5000} \left\{ p \Gamma \frac{da_0}{d\Gamma} + \frac{1}{2} p^2 \Gamma^6 \left( 6a_1 + \Gamma \frac{da_1}{d\Gamma} \right) + \right.$$

$$\left. \frac{1}{5} p^5 \Gamma^{24} \left( 24a_2 + \Gamma \frac{da_2}{d\Gamma} \right) \right\} \quad (C.20)$$

$$\left( s + \frac{R}{10} \ln p \right)_0 = -1.5535 + 1.4286 \ln T + 7.58050 * 10^{-4}T - \frac{46.174}{T} \quad (C.21)$$

Where, s=entropy (KJ/kg K)

## C.9 Superheated Properties for the Region $350^{\circ}C \leq t \leq 500^{\circ}C$

Superheated properties are given by [1].

$$W = \frac{p - RT\rho}{RT\rho^2} = B_8^T(y) * A^T * B_7(x) \quad (C.22)$$

Where, R=4.6152 bar  $cm^3/g \deg K$

p=saturation pressure (bar)

T=saturation temperature (K)

$\rho$  = density ( $g/cm^3$ )

$$A = \begin{bmatrix} -2.933857 & -3.560472 & -0.752570 & -0.355238 & -0.171618 & -0.088253 & -0.023708 & -0.002024 \\ +2.178508 & +1.646874 & +0.257244 & -0.286577 & -0.191781 & -0.118835 & -0.020660 & 0 \\ -0.410626 & -0.712020 & -0.121860 & +0.067088 & +0.080874 & +0.014949 & 0 & 0 \\ +0.027044 & -0.046842 & -0.028628 & -0.032120 & -0.006524 & 0 & 0 & 0 \\ +0.102056 & +0.134922 & +0.052882 & +0.007107 & 0 & 0 & 0 & 0 \\ -0.026332 & -0.049185 & -0.008993 & 0 & 0 & 0 & 0 & 0 \\ +0.002725 & +0.001175 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

$$x = \frac{\rho}{\rho_{max}} \quad (C.23)$$

$$y = \frac{\left(\frac{1}{T} - \frac{1}{T_{max}}\right)}{\left(\frac{1}{T_{min}} - \frac{1}{T_{max}}\right)} \quad (C.24)$$

$$\rho_{max} = 0.70839 \text{ g/cm}^3$$

$$T_{min} = 468.15 \text{ K}$$

$$T_{max} = 1178.15 \text{ K}$$

$$B_m(z) = \begin{bmatrix} b_1(z) \\ b_2(z) \\ \dots \\ \dots \\ b_{m-1}(z) \\ b_m(z) \end{bmatrix} \quad (C.25)$$

$$= \begin{bmatrix} 1 \\ 2z - 1 \\ 2 * b_2 * b_2 - b_1 \\ 2 * b_2 * b_3 - b_2 \\ 2 * b_2 * b_4 - b_3 \\ \dots \\ \dots \\ 2 * b_2 * b_{m-1} - b_{m-2} \end{bmatrix} \quad (C.26)$$

$$\frac{h - h_0}{RT} = \frac{1}{10} \rho_{max} \left( xW + (y + \alpha) \int_0^x \frac{\partial W}{\partial y} dx \right) \quad (C.27)$$

$$h_0 = 1809.25 + 1.48286T + 3.79025 * 10^{-4}T + 46.174 \ln T \quad (C.28)$$

where,  $h$ =enthalpy (kJ/kg)

$$\alpha = \frac{\frac{1}{T_{max}}}{\left( \frac{1}{T_{min}} - \frac{1}{T_{max}} \right)} \quad (C.29)$$

$$\int_0^x \frac{\partial W}{\partial y} dx = B_8^T(x) * C_3 * A * B_8(y) \quad (C.30)$$

$$C_3 = \begin{bmatrix} \frac{1}{2} & -\frac{1}{8} & -\frac{1}{6} & \frac{1}{16} & -\frac{1}{30} & \frac{1}{48} & -\frac{1}{70} \\ \frac{1}{2} & 0 & -\frac{1}{4} & 0 & 0 & 0 & 0 \\ 0 & \frac{1}{8} & 0 & -\frac{1}{8} & 0 & 0 & 0 \\ 0 & 0 & \frac{1}{12} & 0 & -\frac{1}{12} & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{16} & 0 & -\frac{1}{16} & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{20} & 0 & -\frac{1}{20} \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{24} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{28} \end{bmatrix} \quad (C.31)$$

$$\left( s + \frac{R}{10} \ln \rho \right) - \left( s + \frac{R}{10} \ln \rho \right)_0 = -\frac{1}{10} R \rho_{max} \int_0^x W dx + \frac{1}{10} R \rho_{max} (y + \alpha) \int_0^x \frac{\partial W}{\partial y} dx \quad (C.32)$$

$$\left( s + \frac{R}{10} \ln \rho \right)_0 = -2.2593 + 1.02134 \ln T + 7.58050 * 10^{-4} T - \frac{46.174}{T} \quad (C.33)$$

where,  $s$ =entropy (kJ/kg K)

$$\int_0^x W dx = B_8^T(x) * C_3 * A * C_4 * B_7(y) \quad (C.34)$$

$$C_4 = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 8 & 0 & 0 & 0 & 0 & 0 \\ 6 & 0 & 12 & 0 & 0 & 0 & 0 \\ 0 & 16 & 0 & 16 & 0 & 0 & 0 \\ 10 & 0 & 20 & 0 & 20 & 0 & 0 \\ 0 & 24 & 0 & 24 & 0 & 24 & 0 \\ 14 & 0 & 28 & 0 & 28 & 0 & 28 \end{bmatrix} \quad (C.35)$$

values of  $h$ ,  $s$ ,  $v$  are calculated from the equations C.27, C.32, C.23 respectively.

$$h(actual) = h(calculated) - \Delta h \quad (C.36)$$

$$s(actual) = s(calculated) - \Delta s \quad (C.37)$$

$$v(actual) = v(calculated) - \Delta v \quad (C.38)$$

values of  $\Delta h$ ,  $\Delta s$ ,  $\Delta v$  are given in Tables C.1, C.2, C.3 respectively.

Table C.1: Values of  $\Delta h$  ( $350^{\circ}C \leq t \leq 500^{\circ}C$ )

Pressure (bar)	Temperature $t$ ( $^{\circ}C$ )															
	350	360	370	380	390	400	410	420	430	440	450	460	470	480	490	500
20	-5.7	-4.9	-4.1	-3.4	-2.8	-2.2	-1.6	-1.2	-0.8	-0.5	-0.3	-0.2	-0.1	0.0	0.0	0.0
40	-10.3	-9.0	-7.7	-6.5	-5.5	-4.5	-3.6	-2.8	-2.2	-1.7	-1.2	-0.8	-0.5	-0.2	0.0	0.0
60	-13.9	-12.2	-10.6	-9.1	-7.7	-6.4	-5.4	-4.4	-3.5	-2.7	-2.0	-1.4	-0.9	-0.4	-0.1	0.0
80	-16.3	-14.6	-12.8	-11.1	-9.5	-8.1	-6.8	-5.6	-4.5	-3.5	-2.6	-1.9	-1.1	-0.5	-0.2	0.0
100	-17.3	-16.0	-14.3	-12.6	-11.0	-9.5	-8.1	-6.7	-5.4	-4.2	-3.1	-2.2	-1.3	-0.6	-0.2	0.0
120	-15.8	-15.9	-14.9	-13.4	-11.9	-10.5	-9.0	-7.5	-6.1	-4.8	-3.7	-2.6	-1.6	-0.7	-0.2	0.0
140	-10.8	-13.3	-14.0	-13.4	-12.3	-11.1	-9.6	-8.0	-6.5	-5.3	-4.1	-2.9	-1.8	-0.8	-0.2	0.0
160	-0.5	-8.5	-12.0	-12.6	-12.1	-11.1	-9.8	-8.3	-6.9	-5.6	-4.3	-3.1	-1.9	-0.8	-0.2	0.0
180	14.3	3.0	-7.0	-10.7	-11.2	-10.8	-9.7	-8.3	-6.9	-5.6	-4.3	-3.1	-2.0	-0.9	-0.2	0.0
200	13.1	11.0	3.0	-7.0	-9.7	-9.9	-9.3	-8.1	-6.8	-5.5	-4.3	-3.1	-1.9	-0.8	-0.2	0.0
220	12.7	9.2	5.8	-2.9	-7.5	-8.2	-8.3	-7.5	-6.3	-5.1	-4.0	-2.9	-1.8	-0.8	-0.2	0.0

Table C.2: Values of  $\Delta s$  ( $350^{\circ}C \leq t \leq 500^{\circ}C$ )

Pressure (bar)	Temperature $t$ ( $^{\circ}C$ )													
	350	360	370	380	390	400	410	420	430	440	450	460		
20	-0.009	-0.007	-0.006	-0.005	-0.004	-0.003	-0.002	-0.001	-0.001	0.000	0.000	0.000	0	0
40	-0.015	-0.013	-0.011	-0.009	-0.008	-0.006	-0.005	-0.004	-0.003	-0.002	-0.001	-0.001	0.000	0
60	-0.021	-0.018	-0.015	-0.013	-0.011	-0.009	-0.008	-0.006	-0.005	-0.004	-0.003	-0.002	-0.001	0.000
80	-0.024	-0.021	-0.019	-0.016	-0.014	-0.011	-0.010	-0.008	-0.006	-0.005	-0.004	-0.002	-0.001	0
100	-0.025	-0.023	-0.021	-0.018	-0.016	-0.013	-0.011	-0.009	-0.007	-0.006	-0.004	-0.003	-0.002	-0.001
120	-0.023	-0.023	-0.021	-0.019	-0.017	-0.015	-0.012	-0.010	-0.008	-0.006	-0.005	-0.004	-0.002	-0.001
140	-0.015	-0.019	-0.020	-0.019	-0.017	-0.016	-0.013	-0.011	-0.009	-0.007	-0.006	-0.004	-0.002	-0.001
160	0.001	-0.011	-0.017	-0.018	-0.017	-0.016	-0.014	-0.011	-0.009	-0.008	-0.006	-0.004	-0.002	-0.001
180	0.024	0.006	-0.009	-0.015	-0.016	-0.015	-0.014	-0.011	-0.009	-0.008	-0.006	-0.004	-0.003	-0.001
200	0.022	0.018	0.006	-0.009	-0.014	-0.014	-0.013	-0.011	-0.009	-0.008	-0.006	-0.004	-0.002	-0.001
220	0.021	0.015	0.010	-0.003	-0.010	-0.011	-0.012	-0.010	-0.009	-0.007	-0.005	-0.004	-0.002	-0.001

Table C.3: Values of  $\Delta v$  (  $350^{\circ}C \leq t \leq 500^{\circ}C$  )

Pressure (bar)	Temperature $t$ ( $^{\circ}C$ )										
	350	360	370	380	390	400	410	420	430	440	450
20	-0.20	-0.16	-0.13	-0.11	-0.09	-0.07	-0.05	-0.04	-0.03	-0.02	-0.01
40	-0.17	-0.14	-0.12	-0.10	-0.08	-0.06	-0.05	-0.04	-0.03	-0.02	-0.01
60	-0.14	-0.12	-0.10	-0.08	-0.06	-0.05	-0.04	-0.03	-0.02	-0.02	-0.01
80	-0.11	-0.09	-0.08	-0.07	-0.05	-0.04	-0.03	-0.02	-0.02	-0.01	-0.01
100	-0.07	-0.07	-0.06	-0.05	-0.04	-0.04	-0.03	-0.02	-0.02	-0.01	-0.01
120	-0.011	-0.028	-0.035	-0.035	-0.031	-0.027	-0.022	-0.017	-0.013	-0.009	-0.006
140	0.054	0.012	-0.009	-0.016	-0.018	-0.016	-0.014	-0.011	-0.009	-0.006	-0.004
160	0.137	0.063	0.021	0.004	-0.04	-0.006	-0.006	-0.005	-0.004	-0.003	-0.002
180	0.048	0.143	0.063	0.027	0.012	0.005	0.002	0.001	0.001	0.00	0.00
200	0.040	0.050	0.123	0.056	0.032	0.018	0.012	0.008	0.005	0.004	0.002
220	0.033	0.038	0.051	0.057	0.044	0.029	0.020	0.014	0.010	0.006	0.004

## C.10 Superheated Properties for the Region $500^{\circ}C \leq t \leq 800^{\circ}C$

superheated properties for the region  $500^{\circ}C \leq t \leq 800^{\circ}C$  are calculated in a similar way [1] as given in section no. C.9, only difference is that here,

$$\Delta h = 0, \Delta s = 0, \Delta v = 0 \text{ and}$$

$$A = \begin{bmatrix} -2.684145 & -3.110441 & -0.406137 & -0.143284 & -0.063656 & -0.048549 & -0.012639 & -0.000698 \\ +1.571884 & +0.553618 & -0.584342 & -0.801473 & -0.454050 & -0.215289 & 0 & 0 \\ -0.410626 & -0.712020 & -0.121860 & +0.067088 & +0.080874 & +0.014949 & 0 & 0 \\ +0.027044 & -0.046842 & -0.028628 & -0.032120 & -0.006524 & 0 & 0 & 0 \\ +0.102056 & +0.134992 & +0.052882 & +0.007107 & 0 & 0 & 0 & 0 \\ -0.026332 & -0.049185 & -0.008993 & 0 & 0 & 0 & 0 & 0 \\ +0.002775 & +0.001175 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \quad (C.39)$$

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